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<td>Creators</td>
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HD 42659: the only known roAp star in a spectroscopic binary observed with B photometry, TESS, and SALT

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ABSTRACT
We present a multi-instrument analysis of the rapidly oscillating Ap (roAp) star HD 42659. We have obtained B photometric data for this star and use these data, in conjunction with the Transiting Exoplanet Survey Satellite (TESS) observations, to analyse the high-frequency pulsation in detail. We find a triplet that is split by the rotation frequency of the star ($\nu_{\text{rot}} = 0.3756 \text{d}^{-1}; P_{\text{rot}} = 2.66 \text{d}$) and present both distorted dipole and distorted quadrupole mode models. We show that the pulsation frequency, $150.9898 \pm 0.0006 \text{d}^{-1}$ ($P_{\text{puls}} = 9.54 \text{min}$), is greater than the acoustic cut-off frequency. We utilize 27 high-resolution ($R \simeq 65,000$), high signal-to-noise ratio ($\sim 120$) spectra to provide new orbital parameters for this, the only known roAp star to be in a short-period binary ($P_{\text{orb}} = 93.266 \text{d}$). We find the system to be more eccentric than previously thought, with $e = 0.317$, and suggest the companion is a mid-F to early-K star. We find no significant trend in the average pulsation mode amplitude with time, as measured by TESS, implying that the companion does not have an effect on the pulsation in this roAp star. We suggest further photometric observations of this star, and further studies to find more roAp stars in close binaries to characterize how binarity may affect the detection of roAp pulsations.

Key words: asteroseismology – techniques: photometric – techniques: radial velocities – stars: chemically peculiar – stars: individual: HD 42695 – stars: oscillations.

1 INTRODUCTION
The rapidly oscillating Ap (roAp) stars are a rare subclass of the Ap stars that show brightness and radial velocity (RV) variations with periods of 4.7–23.7 min. Discovered using targeted B photometric observations of known Ap stars (Kurtz 1978, 1982), only 76 roAp stars have been identified over a 40-yr period (see catalogues of Smalley et al. 2015; Joshi et al. 2016 and recent works by Balona, Holdsworth & Cunha 2019; Cunha et al. 2019; Hey et al. 2019). With the recent results from the Transiting Exoplanet Survey Satellite (TESS; Ricker et al. 2015), it is becoming evident that the roAp stars are even rarer than previously thought, with an incidence rate among the Ap stars of about 4 per cent (Cunha et al. 2019).

The magnetic Ap stars are characterized by their chemical peculiarities, strong predominately dipolar magnetic fields, and slow rotation. They constitute about 13 per cent of the A stars as a whole (Wolff 1968, 1983). The magnetic fields in the Ap stars are aligned to the rotation axes, and can have magnetic field modulus strengths of up to about 35 kG (Babcock 1960; Elkin et al. 2010), which act to suppress convection, thus allowing gravitational settling of some elements and radiative levitation of others, most notably the rare earth elements. This phenomenon can result in overabundances, in the stellar surface layers, of rare earth elements, which often form spots at the magnetic poles of the Ap stars (Ryabchikova et al. 2007). Those spots are stable over many decades to centuries, thus allowing for the accurate determination of the star’s rotation period through modulation of its light curve, which can be of the order of decades to centuries (Mathys 2015). It is also the magnetic field in the Ap stars that is thought to cause this slow rotation in many of them through magnetic braking (Stepień 2000).

The pulsations in the roAp stars, which are low-degree ($\ell$), high-overtone ($n$) pressure ($p$) modes, are thought to be driven by the $\kappa$-mechanism acting in the HI ionization zone (Balmforth et al. 2001; Cunha 2002). It was found that this mechanism is only viable when the magnetic field suppresses convection around the magnetic poles, explaining why it is the magnetic Ap stars that show these high-overtone pulsation modes that are not found in the non-magnetic A stars. This also provides an explanation as to why the pulsation axis in roAp stars is closely aligned to the magnetic axis. The $\kappa$-mechanism cannot explain the observed pulsation mode frequencies in some stars. In non-magnetic stars there is a theoretical limit to pulsation mode frequencies, the acoustic cut-off frequency, where the wavelength of the pulsation becomes short compared to
Table 1. Basic and derived stellar parameters of HD 42659.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
<th>Reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>Spectral type</td>
<td>Ap SrCrEu</td>
<td>(1)</td>
</tr>
<tr>
<td>$m_V$</td>
<td>$6.79 \pm 0.01$</td>
<td>(2)</td>
</tr>
<tr>
<td>$M_V$</td>
<td>$1.03 \pm 0.23$</td>
<td>(3)</td>
</tr>
<tr>
<td>$\pi$ (mas)</td>
<td>$7.62 \pm 0.04$</td>
<td>(4)</td>
</tr>
<tr>
<td>d (pc)</td>
<td>$131.2 \pm 0.7$</td>
<td>(4)</td>
</tr>
<tr>
<td>$\log T_{\text{eff}}$</td>
<td>$3.90 \pm 0.01$</td>
<td>(3)</td>
</tr>
<tr>
<td>$\log (L/L_\odot)$</td>
<td>$1.48 \pm 0.09$</td>
<td>(3)</td>
</tr>
<tr>
<td>$v \sin i_{\text{mac}}$ (km s$^{-1}$)</td>
<td>$19 \pm 5$</td>
<td>(3)</td>
</tr>
<tr>
<td>$(R_\ast)$ (kG)</td>
<td>$0.40 \pm 0.05$</td>
<td>(3), (6)$^a$</td>
</tr>
<tr>
<td>$M_\ast$ (M$_\odot$)</td>
<td>$2.10 \pm 0.10$</td>
<td>(3)</td>
</tr>
<tr>
<td>$R$ ($R_\ast$)</td>
<td>$2.90 \pm 0.33$</td>
<td>From standard relation</td>
</tr>
</tbody>
</table>


the scale of the surface reflective boundary. However, there are many roAp stars that pulsate above this limit (see fig. 6 of Holdsworth et al. 2018c). It was shown by Cunha et al. (2013) that, in surface regions where convection is not suppressed by the magnetic field, turbulent pressure can be a viable excitation mechanism for the observed supercritical pulsations, i.e. pulsation mode frequencies that are observed above the theoretical acoustic cut-off frequency for a non-magnetic star. In addition, due to the strong magnetic fields in these stars, compressible Alfvén waves are present in the stellar atmosphere, with some, or even most, of the pulsation wave energy residing in the magnetic waves, thus providing the potential for pulsation frequencies higher than the purely acoustic cut-off frequency (Sousa & Cunha 2008).

As previously mentioned, the pulsation axis in roAp stars is misaligned with respect to the rotation axis, but closely aligned to the magnetic one. This can lead to a variation in the viewing aspect of a pulsation over the stellar rotation period. This phenomenon results in amplitude modulation of any pulsation mode with the rotation period, which manifests itself as a multiplet split exactly by the rotation frequency in a Fourier transform. This was first noted by Kurtz (1982) who developed the oblique pulsator model. Subsequently, that model has been revised and improved many times (e.g. Shibahashi & Saio 1985a,b; Shibahashi & Takata 1993; Bigot & Dziembowski 2002; Bigot & Kurtz 2011). This property of the roAp stars gives constraints on the mode geometry, since a multiplet of $2\ell + 1$ components is expected for each mode, where the relative amplitudes of the multiplet components are dependent on the inclination of the rotation axis to the line of sight, $i_{\text{inc}}$, and the obliquity of the magnetic axis with respect to the rotation axis, $\beta$.

The spectroscopic binary fraction of Ap stars is of the order 45 per cent (Carrier et al. 2002; Mathys 2017). However, the short-period ($P < 20$ d) spectroscopic binary fraction amongst the Ap stars is very low, with only about 10 known (Landstreet et al. 2017). Although some roAp stars have been observed to be in visual binary systems (Schöller et al. 2012), there is only one known to be a spectroscopic binary: HD 42659.

2 HD 42659

The subject of this paper is HD 42659 (UV Lep; TIC 33601621); we present basic and derived properties for this star in Table 1.

HD 42659 was discovered to be a pulsationally variable Ap star by Martínez, Kurtz & Ashley (1993). They observed the star through a B filter for a single night with a photometer mounted on the 0.75-m telescope at the South African Astronomical Observatory (SAAO). Later, a more extensive photometric study was conducted by Martínez & Kurtz (1994), where they observed the star for a further 11 nights on the 1.0-m SAAO telescope with the aim using a larger aperture to decrease scintillation noise in the amplitude spectrum. Their results showed variability on only a few nights, indicating strong amplitude modulation of the pulsation, with a period of a few days.

Using a data set combined of 5 nights, which had the highest duty cycle, they concluded that the pulsation had a frequency of $149.9472 \pm 1.7355$ mHz, but stated that this value should be treated with caution due to the mode having low amplitude and the data a complex window pattern. Since those results, no further photometric observations were made.

Spectroscopic observations spanning 4 yr were obtained by Hartmann & Hay (2015) using the High-Accuracy Radial velocity Planet Searcher (HARPS) spectrograph on the 3.6-m telescope at the European Southern Observatory (ESO). Their 48 spectra, obtained on 16 nights, revealed significant RV variations that were successfully modelled as being caused by eccentric binary motion with $P_{\text{orb}} = 93$ d. Although with a period greater than 20 d, this system possesses the potential to study the interplay between the excitation, or suppression, of roAp pulsations by a binary companion.

3 PHOTOMETRY

3.1 The 2017 SAAO observations

We conducted a 3-week observing campaign from SAAO in 2017 December. Using the 1.0-m telescope and one of the Sutherland High Speed Optical Cameras (SHOC; Coppejans et al. 2013), we aimed to obtain the most complete light curve of HD 42659 to date to provide robust values for the pulsation frequency and rotation frequency. The observations were conducted with a B filter, where the signal-to-noise ratio (S/N) is greatest for the roAp stars (Medupe & Kurtz 1998). We provide a log of the observations in Table 2.

The data were reduced using the TEA-PHOT code of Bowman & Holdsworth (2019), which has previously been used for the reduction of SAAO/SHOC data (e.g. Holdsworth et al. 2018a), and binned to 60-s integrations to give each data point the same weighting. Because of the brightness of the star ($V = 6.8$) and the relatively small field of view of the SHOC instrument, there were no suitable comparison stars to perform differential photometry. We therefore removed data that were affected by cloud, then filtered low-frequency variations by pre-whitening the data. This process removed both airmass and sky transparency variations, but also any information on the rotation period of the star. The final light curve is shown in Fig. 1.

In search of the pulsation in HD 42659, we calculate an amplitude spectrum to $300 \text{ day}^{-1}$ and show this in the top panel of Fig. 2. Despite efforts to remove the low-frequency artefacts in the data, the amplitude spectrum shows there is still a significant amount of red noise present. However, at higher frequency where the noise decreases, there is a clear feature that is significant above the noise. A detailed look (bottom panel of Fig. 2) shows that there are two peaks in the amplitude spectrum ($\nu_a = 150.971 \pm 0.007 \text{ day}^{-1}$, $\nu_b = 151.986 \pm 0.008 \text{ day}^{-1}$), separated by 1 day$^{-1}$, which are in the same frequency range as the peak identified in the discovery data (Martínez et al. 1993). With the 2017–2018 SAAO data alone, it is not possible to distinguish which of the peaks represent the
Table 2. Log of 2017/2018 SAAO observations of HD 42659. UTC date and BJD are given for the mid-point of the first exposure. All observations were conducted with the SAAO 1.0-m telescope by DLH (first author).

<table>
<thead>
<tr>
<th>Year</th>
<th>UTC date</th>
<th>BJD (245 0000.0+ )</th>
<th>Duration (min)</th>
</tr>
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<tr>
<td>2017</td>
<td>Dec 12/13</td>
<td>8100.35792</td>
<td>260.2</td>
</tr>
<tr>
<td></td>
<td>Dec 13/14</td>
<td>8101.36511</td>
<td>48.3</td>
</tr>
<tr>
<td></td>
<td>Dec 14/15</td>
<td>8102.35343</td>
<td>381.9</td>
</tr>
<tr>
<td></td>
<td>Dec 15/16</td>
<td>8103.35379</td>
<td>341.9</td>
</tr>
<tr>
<td></td>
<td>Dec 16/17</td>
<td>8104.38644</td>
<td>219.3</td>
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<td>Dec 19/20</td>
<td>8107.41714</td>
<td>273.5</td>
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<tr>
<td></td>
<td>Dec 20/21</td>
<td>8108.47363</td>
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<td></td>
<td>Dec 22/23</td>
<td>8110.32456</td>
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<tr>
<td></td>
<td>Dec 23/24</td>
<td>8111.31963</td>
<td>340.4</td>
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<td></td>
<td>Dec 24/25</td>
<td>8112.31763</td>
<td>353.1</td>
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<td></td>
<td>Dec 25/26</td>
<td>8113.31867</td>
<td>396.2</td>
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<td></td>
<td>Dec 26/27</td>
<td>8114.30889</td>
<td>409.4</td>
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<td></td>
<td>Dec 27/28</td>
<td>8115.31138</td>
<td>395.4</td>
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<td></td>
<td>Dec 28/29</td>
<td>8116.30767</td>
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<td>Dec 30/31</td>
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<td></td>
<td>Dec 31/Jan 01</td>
<td>8119.30476</td>
<td>152.4</td>
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<td>2018</td>
<td>Jan 01/02</td>
<td>8120.28947</td>
<td>457.5</td>
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<tr>
<td></td>
<td>Jan 02/03</td>
<td>8121.29443</td>
<td>448.7</td>
</tr>
<tr>
<td>Total</td>
<td></td>
<td></td>
<td>5693.4</td>
</tr>
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</table>

Figure 1. Reduced 2017–2018 SAAO light curve for HD 42659. Data affected by cloud have been removed, and pre-whitening has been performed to remove airmass and sky transparency variations.

3.2 TESS observations

By far the best photometric data set for HD 42659 was obtained by the TESS mission during sector 6. During its 2-yr mission, TESS will survey almost the entire sky with a cadence of 30 min (Ricker et al. 2015). The spacecraft consists of four cameras in a configuration that images a strip of sky $24^\circ \times 96^\circ$ in size. One camera is centred on the ecliptic pole, while the others extend nearly to the ecliptic plane. The observing strategy consists of 13 sectors per hemisphere that are rotated through every $\sim 27$ d. As such, objects near the ecliptic pole will be observed for about 350 d with objects near the ecliptic plane having just 27 d of data acquired (Ricker et al. 2015).

In addition to the 30-min data, 20 000 stars per sector were chosen to be observed in the high-cadence mode, with observations every 2 min. HD 42659 was included in the 2-min cadence sample along with nearly 1400 other chemically peculiar stars in the context of the roAp programme of the TESS Asteroseismic Science Consortium (TASC). It is those high-cadence data that we use here; they are publicly available and obtainable from the Mikulski Archive for Space Telescopes (MAST) website.

The data for HD 42659 cover a time span of 21.77 d, with a 1.09 d gap for data download. For the analysis, we used the Pre-search Data Conditioned Simple Aperture Photometry (PDCSAP) fluxes, which were produced by the Science Processing Operations Centre (SPOC; Jenkins et al. 2016). The data were converted to magnitudes, and time stamps adjusted for the zero-point correction. The data were of good quality, thus we did not remove any outlying points. We did, however, remove low-frequency instrumental artefacts that remained after the pipeline reductions (a low-frequency trend and some residual background flux that were removed by fitting a series of sine curves). An initial analysis of these TESS data was presented by Balona et al. (2019), but we provide a more detailed analysis here.

At first visual inspection, there is a clear modulation in the light curve of HD 42659 (Fig. 3). This is expected for the Ap stars, which host spots of increased chemical abundances. To determine the rotation period of the star, we calculated the amplitude spectrum using a discrete Fourier transform (Kurtz 1985) and fitted a harmonic series by non-linear least-squares to the data. A total of six frequencies were used such that the amplitude of the highest

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1https://archive.stsci.edu
frequency is at approximately the noise level in the data. We derived a rotation frequency of $0.375566 \pm 0.000035 \, \text{d}^{-1}$, which corresponds to a period of $2.66265 \pm 0.00025 \, \text{d}$. We show a phase folded light-curve plot in Fig. 4.

To analyse the pulsation in the TESS data, we pre-whitened the data in the range $0–20 \, \text{d}^{-1}$ ($0.1–0.23 \, \text{mHz}$) to an amplitude level that is approximately that of the noise at high frequency (i.e. the average peak height in the range $200–300 \, \text{d}^{-1}$ as shown in Fig. 5). Before performing this step, we checked that there is no other signal in this range that could be attributed to an astrophysical signal. After removing the rotation harmonic series, there are only low-amplitude instrumental artefacts remaining.

The pulsation is clearly evident in the top panel of Fig. 5, with a more detailed view in the bottom panel. The multiplet is indicative of a dipole mode, as under the oblique pulsator model, one expects a multiplet of $2\ell + 1$ peaks for a given mode. These high precision TESS data allow us to confirm that the pulsation frequency in this star is $150.9898 \pm 0.0039 \, \text{d}^{-1}$ ($1.747567 \pm 0.000034 \, \text{mHz}$), which is consistent with the highest peak found in the 2017 SAAO data set.

To analyse the multiplet, we iteratively fitted the components non-linear least-squares. We show the results of this fitting procedure in Table 3. After removing the multiplet from the data, there are no further significant peaks remaining. We used this non-linear fit to show that the splitting is equal to the rotation frequency by taking the ratio of the sidelobe splitting to the rotation frequency presented above. For the $\nu - \nu_{\text{int}}$ sidelobe we found the ratio to be $1.004 \pm 0.011$ and for the $\nu + \nu_{\text{int}}$ the ratio was $1.008 \pm 0.013$, both in agreement with the rotation frequency derived above.

Returning to the oblique pulsator model, for a pure, non-distorted mode, one expects, for the right choice of time zero-point, the phases of all components of the frequency multiplet to be equal; this is a signature of pure amplitude modulation. To test this, we assumed that the sidelobes to the pulsation frequency are split from that frequency by exactly the rotation frequency of the star. Therefore, using the rotation frequency derived above, we force fitted the sidelobes, then fitted the light curve by linear least-squares, choosing a zero-point in time such that the phases of the sidelobes are equal. The results of this test are shown in Table 4.

It is evident that the phases are in excellent agreement, showing that the multiplet is a result of pure amplitude modulation, a signature of oblique pulsation. The amplitudes of the sidelobes, although unequal as a result of the Coriolis force (e.g. Shibahashi & Takata 1993; Bigot & Dziembowski 2002), are in agreement within the errors. It seems that HD 42659 is a good example of a ‘well-behaved’
roAp star. This makes the application of the full oblique pulsator model to this star valid.

In the absence of limb darkening and spots, Kurtz, Shibashashi & Goode (1990) provided a relationship among the amplitudes of a dipole triplet, and the geometry of the star:

\[ \tan \theta \tan \beta = \frac{A_{\pm1}^{i} + A_{0}^{i}}{A_{0}^{i}}, \]  

(1)

where \(A_{\pm1}^{i}\) are the sidelobe amplitudes, \(A_{0}\) is the amplitude of the pulsation mode, \(\theta\) and \(\beta\) are the angles of inclination and magnetic obliquity, respectively. Using the results from Table 3 we find that \(\tan \theta \tan \beta \approx 1.68 \pm 0.26\). Although, without direct information on either \(\theta\) or \(\beta\), we are only able to provide constrains on the two angles, which are shown in Fig. 6.

We show the sum of \(\theta\) and \(\beta\) values as these enable visualization of the star. For a pure dipole mode, the equator is a node and both poles are antinodes. Given that we see both magnetic poles, as suggested by the double-wave nature of the rotational variation of the light curve, it is likely we see both pulsation poles. For this to be the case, \(\theta + \beta\) must be >90°, which is satisfied by the results of equation (1) and Fig. 6. In this case, the equatorial line crosses the line of sight and the amplitude should go to zero and a phase change of \(\pi\) rad is expected.

To test this, we split the light curve into discrete segments of 0.5 d and calculated the amplitude and phase at fixed frequency, with the results shown in Fig. 7. Although the scatter is large, due to the low mode amplitude, there is an obvious change in the pulsation amplitude and phase over the rotation period. This is consistent with the oblique pulsator model.

In conducting this test, we noticed an interesting phenomenon. Although the choice of segment length is arbitrary, as long as there is sufficient frequency resolution, we often choose integer multiples of the pulsation period. We did this and found, when phasing the data on the rotation period, our amplitude and phase measurements all fell at discrete rotation phases. Further investigation showed that the rotation period is \(401.94 \pm 0.06\) times the pulsation period, i.e. an integer multiple within the errors. Whether this is coincidence or related is yet to be determined.

Furthermore, the high scatter in the amplitude plot in Fig. 7 aroused suspicion, as one would expect the pulsation amplitude to be similar at the same phase for different rotation cycles (cf. fig. 15 in Holdsworth et al. 2016). To investigate this further, we calculated the pulsation amplitude at a fixed frequency in segments of the same length of the rotation period. In this way, we sample the same average pulsation amplitude over one complete rotation to the next. The results, shown in Fig. 8, suggest that the average pulsation amplitude may not be stable over the 21-d TESS observations. We therefore fitted a linear regression to the data and found a gradient of \(-1.0 \pm 0.7\) μmag d\(^{-1}\), with a \(p\)-value of 0.32 and a \(\chi^{2}\) of 2.26. This fit showed that there is no statistically significant change to the pulsation amplitude over the TESS observations period. We attribute the scatter in Fig. 7 to the low S/N of the data.

Our final task with the TESS data is to model the amplitude and phase variations in Fig. 7 following the method of Saio (2005), as employed and discussed by Holdsworth et al. (2016). The magnetically distorted eigenfunction of a roAp pulsation mode is expanded using axisymmetric spherical harmonics, \(Y_{\ell}^{m}\). The contribution of each \(\ell\) component to the observed light variations is...
Figure 9. Left: best-fitting distorted dipole model to the observed amplitude and phase variations. We show a representative model, as any values of $i_{\text{inc}}$, $\beta$ that satisfy equation (1) and Fig. 6 provide good fits to the observations. Right: the best-fitting distorted quadrupole model. In this case, only a small value of the polar magnetic field strength allows us to obtain a good fit as stronger fields predict smaller phase variations due to a more significant contribution from the spherically symmetric component. In both cases, the red colour represents the observations, with the black being the models. The dotted lines represent the non-distorted case in each model.

Unfortunately, the intrinsic low amplitude of the pulsation mode, coupled with the red TESS filter, means that the observations do not have sufficient precision to allow us to conclude confidently on the degree of the mode in HD 42659; however, the evidence we have presented leads us to favour a dipole mode for the pulsation in HD 42659.

4 SPECTROSCOPY

Since HD 42659 is the only roAp star known to be in a spectroscopic binary system (Hartmann & Hatzes 2015), it provides us with the opportunity to test whether the binary companion affects the pulsations in the Ap star. To that end, we obtained 27 spectra with the High-Resolution Spectrograph (HRS; Bramall et al. 2010; Crause et al. 2014) mounted on the Southern African Large Telescope (SALT; Buckley, Swart & Meiring 2006). Observations were taken in High Resolution (HR) mode, $R \approx 65,000$, with an integration time of 158 s. The S/N varied between about 100 and 150 in the central region of the spectrum, dependent on observing conditions.

HRS is a dual-beam spectrograph with wavelength coverage of 3700–5500 and 5500–8900 Å. In the following analysis, we chose to use only the blue arm data. This removed any issues with telluric line contamination (which is significant given the elevation of the SAAO site) and removed any systematic offsets between the blue and red arm wavelength calibrations.

The observations were automatically reduced using the SALT custom pipeline, which is based on the ESO’s MIDAS pipeline (Kniazev, Gvaramadze & Berdnikov 2016, 2017). The data underwent standard calibrations, including flat-fielding, order extraction, and wavelength calibration. Finally, the orders were merged to provide a 1D continuous spectrum for each arm. Finally, we applied a barycentric velocity correction to the observations.
the light curve to ensure we had not pre-whitened or neglected any of the secondary star in the spectra that, when considering the shape of the CCF.

The final RV measurements are shown in Table 5.

Normalizing the spectra. The final RV measurements are shown in Fig. 10 and Table 5.

normalizing the spectra. The final RV measurements are shown in Fig. 10 and Table 5.

We fitted the final RV measurements with the RVLIN code of Wright & Howard (2009), and included the bootstrapping method for error analysis by Wang et al. (2012). We provided initial estimates of the system parameters based on those presented by Hartmann & Hatzes (2015) and the mass estimate from Kochukhov & Bagnulo (2006). The results of the fitting process are presented by Hartmann & Hatzes (2015), with only the period we derive being similar. Our solution suggests a much more eccentric orbit (i.e., with a much larger velocity semi-amplitude). When inspecting Fig. 3 of Hartmann & Hatzes (2015), it is noted that their orbital phase coverage is not complete at the very bottom of the velocity curve that perhaps has led to these discrepancies. The difference in these two parameters has an effect on the minimum mass of the companion, which in our case is larger at 0.69 ± 0.01 M⊙. We see no indication of the secondary star in the spectra that, when considering the mass function and magnitude ratios, implies that the companion is between mid-F and early-K. In light of this, we double checked the light curve to ensure we had not pre-whitened or neglected any other signal that could be attributed to the companion. None was found.

With the binary solution, we were able to predict the epochs of eclipse events using the EXOFAST routine of Eastman, Gaudi & Agol (2013). The TESS observations coincide with one possible epoch, as denoted by the vertical dashed line in Fig. 10. However, we see no evidence of an eclipse in the light curve.

4.1 Radial velocity measurements

Although the SALT pipeline provides heliocentric radial velocities (RVs) of a given spectrum, we chose to perform our own RV analysis. To do this, we first synthesized a spectrum using ISPEC (Blanco-Cuaresma et al. 2014; Blanco-Cuaresma 2019) using the stellar parameters in Table 1, then used this synthetic spectrum to normalize the SALT observations, also using ISPEC.

To derive the RV measurements, we used the cross-correlation function (CCF) technique. Each observation was cross-correlated with the synthetic spectrum, with the RV being taken when CCF is maximized. The RV, and error, was derived by a Gaussian fit to the shape of the CCF.

Finally, with an RV measurement for each spectrum, we repeated the above process but removed the initial RV, then reapplied it, during the normalization process. This served to mitigate any errors introduced to the initial RV measurement when automatically normalizing the spectra. The final RV measurements are shown in Fig. 10 and Table 5.

We fitted the final RV measurements with the RVLIN code of Wright & Howard (2009), and included the bootstrapping method for error analysis by Wang et al. (2012). We provided initial estimates of the system parameters based on those presented by Hartmann & Hatzes (2015) and the mass estimate from Kochukhov & Bagnulo (2006). The results of the fitting process are shown in Table 6.

We find significantly different results for most parameters to those presented by Hartmann & Hatzes (2015), with only the period we derive being similar. Our solution suggests a much more eccentric orbit (e = 0.317 ± 0.004 compared to e = 0.146 ± 0.027), and a much larger velocity semi-amplitude. When inspecting Fig. 3 of Hartmann & Hatzes (2015), it is noted that their orbital phase coverage is not complete at the very bottom of the velocity curve that perhaps has led to these discrepancies. The difference in these two parameters has an effect on the minimum mass of the companion, which in our case is larger at 0.69 ± 0.01 M⊙. We see no indication of the secondary star in the spectra that, when considering the mass function and magnitude ratios, implies that the companion is between mid-F and early-K. In light of this, we double checked the light curve to ensure we had not pre-whitened or neglected any other signal that could be attributed to the companion. None was found.

With the binary solution, we were able to predict the epochs of eclipse events using the EXOFAST routine of Eastman, Gaudi & Agol (2013). The TESS observations coincide with one possible epoch, as denoted by the vertical dashed line in Fig. 10. However, we see no evidence of an eclipse in the light curve.

5 DISCUSSION AND CONCLUSIONS

The ground-based photometric observations of HD 42659 do not, unfortunately, provide significantly more information than was previously known. However, the TESS and spectroscopic data do add to our understanding of the roAp star, and the system as a whole.

The TESS photometric observations are the first data that allow us to measure the rotation period of this star. First done by Balona et al. (2019), we have refined the fit here to arrive at a rotation period of 2.66265 ± 0.00025 d. This finding allows us to shed some

![Figure 10. SALT/HRS RV measurements and derived fit. The blue shading around the start of 2018 indicates the time period that SAAO photometry was obtained, with the orange shading around the start of 2019 shows the TESS data range. The vertical dashed line in the orange shading represents a time when an eclipse is expected, if the inclination allows.](https://academic.oup.com/mnras/article-abstract/489/3/4063/5558256 by University of Central Lancashire user on 18 September 2019)
light on the number of nights that Martinez & Kurtz (1994) failed to measure a pulsation; their data, as shown in their table 5 and figure, are separated such that pulsation maximum was often not observed, thus the pulsation was below their detection limit. It was fortunate that they persevered and obtained good rotation phase coverage of the star, although they did not know what that was at the time.

The TESS data have also allowed us to provide the first models of the pulsation in HD 42659. We have presented two model fits to the data and have a preference for the dipole fit as the quality of the fit is better (with a reduced $\chi^2$ of 4.2 rather than 5.8 for the quadrupole model). However, it is only when considering a distorted quadrupole mode that we can provide estimates for the geometry and polar magnetic field strength in the star. Using the stellar rotation period derived here, we can estimate an equatorial rotation velocity of $v_{\text{eq}} = 55 \pm 6$ km s$^{-1}$ using the value of the radius, calculated with the standard equation, in Table 1. This result, and knowledge of the $\sin i_{\text{rot}}$ from Elkin et al. (2008), allows us to independently estimate $i_{\text{rot}}$ at $20^{\circ} \pm 3^{\circ}$. This value is in agreement with the pulsation modelling results presented earlier, and is also in agreement with the results of equation (1) when considering a dipole mode. We highlight, in Fig. 6, the estimated values of $i_{\text{rot}}$ from this discussion.

We do not assume that the orbital axis is aligned with rotation axis of the Ap star. If our estimate of $i_{\text{rot}} = 20^{\circ} \pm 3^{\circ}$ is correct, and if the orbital and rotation axes are aligned (i.e. $i_{\text{rot}} = i_{\text{orb}}$), this leads to a secondary mass range of $M_2 \gtrsim 2.4-4.0$ M$_{\odot}$. Since we see no sign of a secondary component in the spectra (see end of Section 4), an orbital inclination angle of $\gtrsim 35^{\circ}$ is required, giving a secondary mass range of $M_2 \approx 0.7-1.4$ M$_{\odot}$.

As discussed earlier, this star has previously been identified as a possible distorted quadrupole pulsator from its position in the $T_{\text{eff}}-\varpi L/\varpi M$ plane (see Saio 2014, for details) and its proximity to other roAp stars pulsating with distorted quadrupole modes. All of these stars pulsate above their theoretical acoustic cut-off frequency, $v_{\text{ac}}$. This theoretical limit can be calculated by scaling from the solar value such that

$$v_{\text{ac,\odot}} = \frac{M/M_{\odot}(T_{\text{eff}}/T_{\text{eff,\odot}})^{3.5}}{L/L_{\odot}}, \quad (3)$$

where $v_{\text{ac,\odot}}$ is taken to be $458$ d$^{-1}$ (Jiménez, García & Pallé 2011). Using values from Table 1, we find $v_{\text{ac}}$ is about $100 \pm 10$ d$^{-1}$. This is significantly lower than the measured pulsation mode frequency ($\sim 151$ d$^{-1}$), thus confirming that HD 42659 is a roAp star pulsating at a supercritical frequency. Since linear, non-adiabatic models show that the $k$-mechanism can only excite modes close to, but not beyond, the acoustic cut-off frequency, it is possible that excitation mechanism is not driving the pulsation in HD 42659. However, the models do not directly account for the magnetic field that plays a significant role in these supercritical pulsators. Alternatively, Cunha et al. (2013) showed that models where convection is not suppressed by the magnetic field some supercritical pulsations could be driven. If this is the case here, turbulent pressure could be the driving mechanism for the pulsation in HD 42659.

HD 42659 is currently the only roAp known to be in a relatively short, spectroscopic, orbit. Other roAp stars, such as $\alpha$ Cir and $\gamma$ Equ (Schöller et al. 2012), do have companions, but these are in wide orbits. HD 42659 allows us to directly probe the influence of binarity on the roAp phenomenon, as suggested by Hartmann & Hatzes (2015). The TESS data, and our simultaneous spectra, allow us to do this for the first time for a roAp star. However, there are Ap stars in close binaries that have been studied in this way. Recently, Skarka et al. (2019) studied HD 99458 both spectroscopically and photometrically. They found a primary Ap star showing $\delta$ Sct pulsations (i.e. low-overtone $p$-modes) in an eclipsing binary with an M-dwarf companion. Clearly the $\delta$ Sct pulsations were not suppressed, but the authors did not investigate this to its full potential. Furthermore, TESS observations are uncovering more Ap stars in short-period eclipsing binary systems, of which there is currently a dearth (e.g. Mathys 2017, and references therein), providing more opportunities to investigate these systems.

The TESS data cover just under a quarter of the binary orbit, and occur on the approach to periastron passage. The orange shading in Fig. 10 shows this visually. For a clearer view of this, we present a phase folded RV curve in Fig. 11 and below this we plot the normalized average amplitude of the pulsation, as calculated over a rotation cycle. We plot both the SAAO and TESS results in this way. These amplitudes cannot be directly compared, and are normalized to their respective maximum.

The SAAO amplitude results are inconclusive due the pulsation S/N and the gaps in the data, which mean we are sampling different pulsation amplitudes due to the rotation. The TESS data do show a trend of $-1.0 \pm 0.7$ μmag d$^{-1}$ that is not statistically significant. Therefore, it is not possible to conclude whether the binary companion is affecting the pulsation mode amplitude.

It is unfortunate that the TESS data cover such a small fraction of the orbital phase curve, but this adds further motivation to obtain high-cadence, high-quality photometric observations of HD 42659 at other orbital phases, and especially at apastron where we currently have no photometric data. Because of the low amplitude of the mode, and the need for high-precision data, this target is a good candidate for the Whole Earth Telescope (Sullivan 2001; Provencal et al. 2014). TESS observations are enabling us to find new roAp stars, and reclassify known Ap stars that were thought to be constant from ground-based observations as variable (Balona et al. 2019; Cunha et al. 2019). Those targets are prime candidates to perform long-term spectroscopic studies to search for binary companions that may affect any pulsation mode amplitudes.

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