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






























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# Magnetic Fields toward Ophiuchus-B Derived from SCUBA-2 Polarization Measurements

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## Abstract

We present the results of dust emission polarization measurements of Ophiuchus-B (Oph-B) carried out using the Submillimetre Common-User Bolometer Array 2 (SCUBA-2) camera with its associated polarimeter (POL-2) on the James Clerk Maxwell Telescope in Hawaii. This work is part of the *B*-fields in Star-forming Region Observations survey initiated to understand the role of magnetic fields in star formation for nearby star-forming molecular clouds. We present a first look at the geometry and strength of magnetic fields in Oph-B. The field geometry is traced over  $\sim 0.2$  pc, with clear detection of both of the sub-clumps of Oph-B. The field pattern appears significantly disordered in sub-clump Oph-B1. The field geometry in Oph-B2 is more ordered, with a tendency to be along the major axis of the clump, parallel to the filamentary structure within which it lies. The degree of polarization decreases systematically toward the dense core material in the two sub-clumps. The field lines in the lower density material along the periphery are smoothly joined to the large-scale magnetic fields probed by NIR polarization observations. We estimated a magnetic field strength of  $630 \pm 410 \mu\text{G}$  in the Oph-B2 sub-clump using a Davis–Chandrasekhar–Fermi analysis. With this magnetic field strength, we find a mass-to-flux ratio  $\lambda = 1.6 \pm 1.1$ , which suggests that the Oph-B2 clump is slightly magnetically supercritical.

**Key words:** ISM: individual objects (Ophiuchus) – ISM: structure – polarization – radio continuum: ISM – stars: formation

## 1. Introduction

Low-mass stars are formed in dense cores ( $M \approx 1\text{--}10 M_{\odot}$ , size  $\approx 0.1\text{--}0.4$  pc and density  $\approx 10^4\text{--}10^5 \text{cm}^{-3}$ ) embedded in molecular clouds that are generally self-gravitating, turbulent, magnetized, and thought to be compressible fluids, and are expected to form one or a few stars when they become unstable to gravitational collapse.

Considering the magnetized nature of molecular clouds (Shu et al. 1987; McKee & Ostriker 2007), we expect magnetic fields (*B*-fields) to also have a significant impact on dense cores. Nevertheless, the role of *B*-fields on the formation of dense cores and their evolution into the various stages of star formation is still under debate. Several observational and theoretical studies have been dedicated to understand the importance of *B*-fields in star formation. For example, isolated low-mass cores that are magnetically dominated may gradually condense out of a large-scale cloud through ambipolar diffusion (e.g., Shu et al. 1987; McKee et al. 1993; Mouschovias & Ciolek 1999;

Allen et al. 2003). In this picture, the core will be flattened into a disk-like morphology on scales of a few thousand au, with field lines primarily parallel to the symmetry axis. These field lines become pinched into an hour-glass morphology as mass accumulates in the core and self-gravity becomes more significant (Fiedler & Mouschovias 1993; Galli & Shu 1993; Girart et al. 2006; Attard et al. 2009). On the other hand, *B*-fields are expected to be less significant if cores form via turbulent flows (Mac Low & Klessen 2004; Dib et al. 2007, 2010). In this picture the *B*-field morphology will be more chaotic (Crutcher 2004; Hull et al. 2017b).

*B*-fields are often characterized by dust polarization observations. Dust grains are expected to align with their short axes parallel to an external *B*-field. As a result, thermal dust emission at sub-millimeter or millimeter wavelengths from these grains will be polarized perpendicular to the *B*-field (e.g., Cudlip et al. 1982; Hildebrand et al. 1984; Hildebrand 1988; Rao et al. 1998; Dotson et al. 2000, 2010; Lazarian 2000; Vaillancourt & Matthews 2012; Hull et al. 2014). The actual mechanism by which dust grains align with a *B*-field is still unclear (Lazarian 2007; Andersson et al. 2015). The most accepted mechanism to date is radiative torque alignment (e.g., Hoang & Lazarian 2014; Andersson et al. 2015; Hoang et al. 2015), which was originally proposed by Dolginov and Mitrofanov (1976).

Polarized thermal dust emission is important because it probes the magnetic fields in the denser regions with extinction

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**Table 1**  
Log of SCUBA-2/POL-2 Observations

No. of Observations	Observation Date (year, month, date)
3	2016 Apr 27
3	2016 Apr 28
5	2016 May 02
1	2016 May 05
5	2016 May 07
2	2016 May 10
2	2016 May 11

$A_V > 50$  mag, since NIR/optical polarization measurements are limited by the number of detectable background stars and cannot probe these high-density environments. At sub-mm wavelengths, one can probe the  $B$ -field morphology deep inside the dense cores ( $n_{\text{H}_2} \sim 10^5\text{--}10^6 \text{ cm}^{-3}$ ) where the central protostars and their circumstellar disks form.  $B$ -fields in more diffuse medium ( $A_V \sim 1\text{--}20$  mag), associated with larger-scale structures ( $>1$  pc), can typically be probed using dust extinction polarization of background stars at optical or NIR wavelengths (Hall 1949; Hiltner 1949a, 1949b).

Ophiuchus is a molecular cloud at a distance of  $\approx 140$  pc (e.g., Chini 1981; de Geus et al. 1989; Knude & Hog 1998; Rebull et al. 2004; Loinard et al. 2008; Ortiz-León et al. 2017). Therefore this is one of the closest low-mass star-forming regions (Wilkings 1992; Johnstone et al. 2000; Padgett et al. 2008; Miville-Deschênes et al. 2010). The Ophiuchus region is highly structured, with several  $\sim 1$  pc sized clumps. This study focuses on the Oph-B clump, which has two active star-forming sub-regions named Oph-B1 and Oph-B2, separated by  $\sim 5'$ . The Oph-B2 clump is the second densest region in Ophiuchus.

We present the first sub-mm polarization observations made using the Submillimetre Common-User Bolometer Array 2 (SCUBA-2) camera with the POL-2 polarimeter toward the Oph-B region to map the  $B$ -fields on core scales. The molecular line observations toward this region using the Heterodyne Array Receiver Program are already available in White et al. (2015) to understand the kinematics of the cloud.

The paper is organized as follows. In Section 2, we describe the observations and data reduction; in Section 3, we give initial results; in Section 4, we analyze the polarization data and estimate the magnetic field strength; and in Section 5, we summarize the paper.

## 2. Data Acquisition and Reduction Techniques

We observed Oph-B in  $850 \mu\text{m}$  polarized emission with SCUBA-2 (Holland et al. 2013) in conjunction with POL-2 (Friberg et al. 2016; P. Bastien et al. 2018, in preparation) as part of the  $B$ -fields in Star-forming Region Observations (Ward-Thompson et al. 2017) survey under project code M16AL004 at the James Clerk Maxwell Telescope (JCMT). This survey aims to use dust polarization maps of nearby molecular clouds to probe  $B$ -field structures. Oph-B was observed over a series of 21 observations, with an average integration time of  $\sim 0.55$  hours per observation, in good weather ( $0.05 < \tau_{225} < 0.08$ , where  $\tau_{225}$  is atmospheric opacity at 225 GHz). Table 1 summarizes the observation logs. The scan pattern used was a POL-2 daisy (Friberg et al. 2016) producing a uniform, high signal-to-noise coverage over the central  $3'$  of the map. This observing mode is based on the

SCUBA-2 CV daisy scan pattern (Holland et al. 2013) but modified to have a slower scan speed ( $8''/\text{s}$  compared to  $155''/\text{s}$ ) to obtain sufficient on-sky data for good Stokes  $Q$  and  $U$  values. This pattern maps a fully sampled,  $12'$ -diameter, circular region with an effective resolution of  $14''.1$ . Coverage decreases, with a consequent significant increase in the rms noise, toward the edges of the map. The wave plate rotates at a frequency of 2 Hz.

The data were reduced using the `pol2map` python script in the Starlink (Currie et al. 2014) SMURF (Jenness et al. 2013) package. The reduction followed three main steps. In step one, `pol2map` uses the `calccu` command to create Stokes  $Q$ ,  $U$ , and  $I$  timestreams from the raw data. It then creates an initial Stokes  $I$  map for each of the observations using the `makemap` (Chapin et al. 2013) routine, and coadds these to create an initial estimate of the Stokes  $I$  emission in the region. In step two, `pol2map` is re-run. The initial  $I$  map of the region is used to generate a fixed signal-to-noise-based mask, which is used to define regions of astrophysical emission and thus to consistently re-reduce the previously created Stokes  $I$  timestreams for each observation (see Chapin et al. 2013 for a detailed discussion on `makemap` and Mairs et al. 2015 for the role of masking in SCUBA-2 data reduction). The new set of Stokes  $I$  maps are then coadded to produce the final Stokes  $I$  emission map of the region. In step 3, `pol2map` is again re-run, this time creating a Stokes  $Q$  and  $U$  maps for each observation from their timestreams, using the same mask as used in step 2, and coadding these to produce final Stokes  $Q$  and  $U$  emission maps for the region. The final Stokes  $Q$ ,  $U$ , and  $I$  maps, and their corresponding variance maps, are then used to create a polarization vector catalog. The vectors are debiased to remove the effect of statistical biasing in low signal-to-noise-ratio (S/N) regions. See Mairs et al. (2015) and Pattle et al. (2017) for a detailed description of the SCUBA-2 and POL-2 data reduction process, respectively.

The rms noise in the maps was estimated by selecting an emission-free region near the map center and finding the standard deviation of the measured flux density distribution in that region. A rms noise value of  $\sim 2$  mJy/beam (Pattle et al. 2017) was set as the target value for the BISTRO survey. The estimated noise in the Stokes  $I$  map of Oph-B is  $\sim 3.5$  mJy/beam.

Bias in the measured polarization fraction and polarized intensity values result from the fact that polarized intensity is defined as being positive, and so uncertainties on Stokes  $Q$  and  $U$  values tend to increase the measured polarization values (Vaillancourt 2006; Kwon et al. 2018). Debiased values of polarized intensity are estimated to be

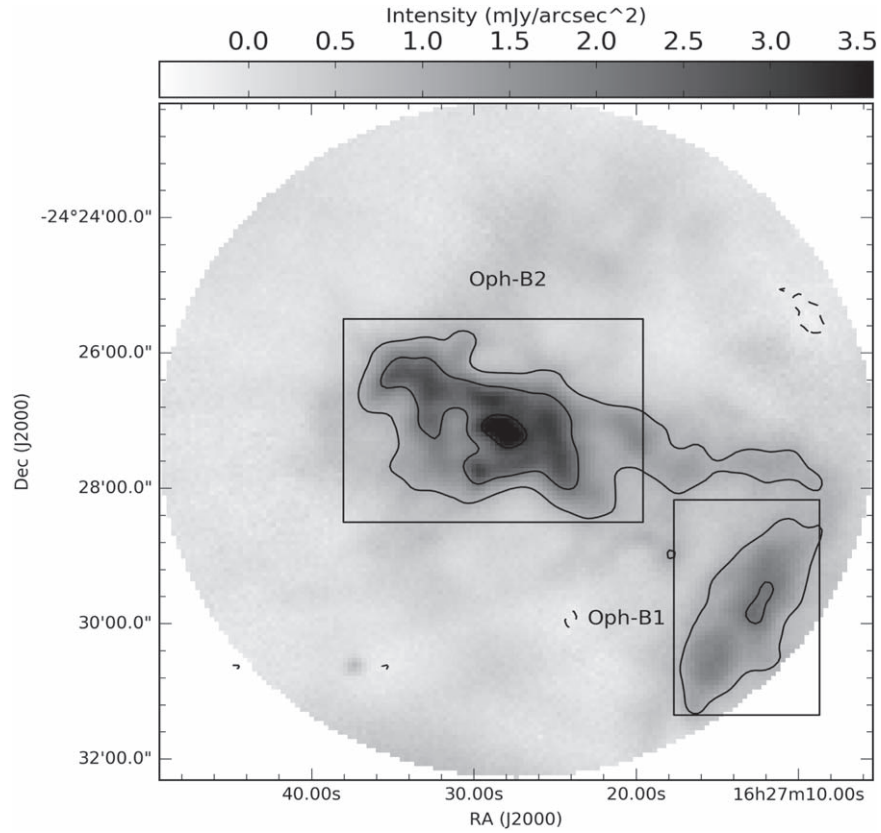
$$\text{PI} = \sqrt{Q^2 + U^2 - 0.5(\delta Q^2 + \delta U^2)}, \quad (1)$$

where PI is polarized intensity,  $\delta Q$  is the uncertainty on Stokes  $Q$ , and  $\delta U$  is the uncertainty on Stokes  $U$ . The polarization fraction,  $P$ , is then given by

$$P = \frac{\text{PI}}{I}, \quad (2)$$

where  $I$  is the total intensity.

The polarization position angle  $\theta$  and the corresponding uncertainties (Naghizadeh-Khouei & Clarke 1993) are



**Figure 1.** 850  $\mu\text{m}$  dust emission map of Oph-B clump obtained using SCUBA-2 on the JCMT. The two sub-clumps, B1 and B2, are labeled with boxes. Contours are of 850  $\mu\text{m}$  emission (grayscale) and the contour levels are drawn with (0.6, 1.8, 2.9, 4.1) mJy/arcsec<sup>2</sup>.

estimated as

$$\theta = \frac{1}{2} \tan^{-1} \left( \frac{U}{Q} \right), \quad (3)$$

and

$$\delta\theta = 0.5 \times \frac{\sqrt{(Q^2 \times \delta U^2 + U^2 \times \delta Q^2)}}{(Q^2 + U^2)} \times \frac{180^\circ}{\pi}. \quad (4)$$

The debiased polarized intensity, polarization fraction, and angle derived as explained above are used in the analysis of the paper. Magnetic field orientation is derived by rotating the  $\theta$  by  $90^\circ$ . The details of the procedure followed for producing polarization values are explained in detail by Kwon et al. (2018).

### 3. Results

Figure 1 shows the Stokes  $I$  map of Oph-B obtained from SCUBA-2/POL-2 observations. The Oph-B clump contains two structures, one in the southwest known as Oph-B1 and other in the northeast known as Oph-B2 (Motte et al. 1998). These structures are labeled with boxes in Figure 1. The entire clump appears to be rather cold, with dust temperatures ranging from  $\sim 7$  to 23 K at the center of the clump (Stamatellos et al. 2007).

The Oph-B region contains at least three Class I protostars and five flat spectrum objects, many of which are driving outflows (White et al. 2015). Figure 2 shows the distributions of young stellar objects (YSOs) in Oph-B with the blue- and redshifted lobes of the molecular outflows driven by IRS47, as

identified by White et al. (2015), marked. Figure 2 also shows dense condensations identified using Gaussclump (Stutzki & Guesten 1990) in  $\text{N}_2\text{H}^+$  (1-0) data by André et al. (2007).

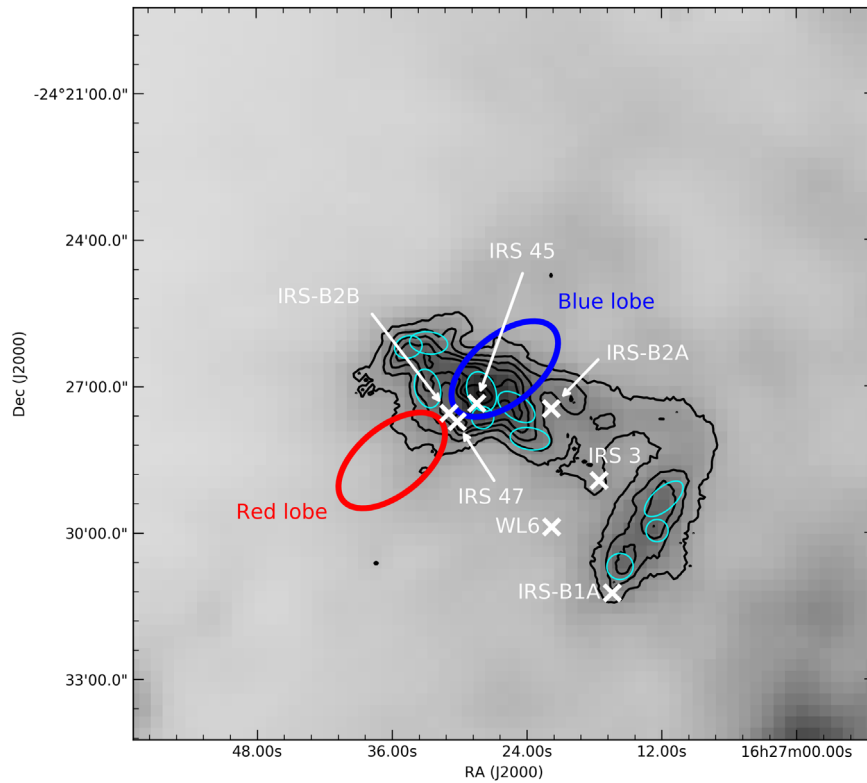
$B$ -fields derived from dust polarization measurements are shown by white line segments in the left-hand panel of Figure 3. The polarization vectors are rotated by  $90^\circ$  to show the  $B$ -field orientation. The length of the vectors shows the polarization fraction. The underlying image is the Stokes  $I$  map, showing the geometry of the cloud. For the following analysis we used those polarization measurements where  $\frac{\text{PI}}{\delta\text{PI}} > 2$  in order to ensure we have a large and robust sample. The measurements with  $\frac{\text{PI}}{\delta\text{PI}} > 3$  are also shown as cyan color vectors in Figure 3. PI and  $\delta\text{PI}$  are the polarized intensity and uncertainty on polarized intensity, respectively. The  $B$ -field orientations obtained from NIR polarization measurements by Kwon et al. (2015) are shown in the figure as red line segments. The right-hand panel of Figure 3 shows the  $B$ -field orientation, with polarization vectors normalized. Here the lengths of vectors are independent of the polarization fraction.

Figure 4 shows the polarization vectors (rotated by  $90^\circ$ ) with the data averaged over 4 pixels using  $2 \times 2$  binning, showing the general morphology of the  $B$ -fields.

Figure 5 compares the distribution of position angle in Oph-B1 and Oph-B2 using the POL-2 data from Figure 3. The figure shows histograms of each complex and of the entire Oph-B clump.

#### 3.1. Comparison of SCUBA-POL and SCUBA-2/POL-2 Results

Matthews et al. (2009) cataloged polarization observations for various star-forming regions using 850  $\mu\text{m}$  data obtained



**Figure 2.** This figure shows sub-mm continuum emission and outflow-related structures and YSO locations in Oph-B. The background image shows Herschel 500  $\mu\text{m}$  emission. Contours show the 850  $\mu\text{m}$  SCUBA-2 dust continuum emission. The positions of YSOs identified in the region are marked with white crosses and labeled. Blue and red ellipses show the blue and redshifted wings of the prominent outflow associated with IRS47 (White et al. 2015), respectively. Cyan ellipses show the dense structures identified in  $\text{N}_2\text{H}^+$  data (André et al. 2007). IRS-B1A, IRS-B2A, and IRS-B2B correspond to YSOs SSTc2dJ162716.4-243114, SSTc2d162721.8-242728, and SSTc2d162730.9-242733 in the SIMBAD database, respectively.

using SCUBA-POL (SCUPOL) at the JCMT. Figure 6 compares the polarization vectors (not rotated by  $90^\circ$ ) toward Oph-B2 (Oph-B1 was not observed with SCUPOL). We found 27 SCUPOL data points with coordinates that matched those of our observations within  $10''$ . The polarization vectors in common agree well in position angles. Figure 7 shows a comparison of degree of polarization and position angle between the two data sets. The polarization angles measured by the two instruments agree with each other within the errors (Figure 7, upper panel). In contrast, the polarization fractions measured with SCUPOL are systematically higher than those with POL-2, especially in the low  $P$  values data (Figure 7, lower panel), suggesting that such a difference is caused by the fact that errors in polarization fraction,  $P$ , do not obey Gaussian statistics, since  $P$  can never be negative and hence, as discussed above, absolute values of  $P$  observed with all instruments and at all wavelengths tend to increase in areas of very low S/N. This has been discussed in detail by Vaillancourt (2006), placing the confidence level in polarization measurements. The SCUPOL data have lower S/N than our own data, and consequently the polarization fractions measured with that instrument are typically higher than those measured with POL-2.

We also performed the Kolmogorov–Smirnov (K–S; Press et al. 1992) statistical test to quantify the deviation between SCUPOL and POL-2 polarization position angles toward Oph-B2 and found a probability of 0.905. The average difference in the orientation between the position angles of SCUPOL and POL-2 segments is found to be  $\sim 20^\circ$ . The distribution of the differences in the orientations of position angles in these two

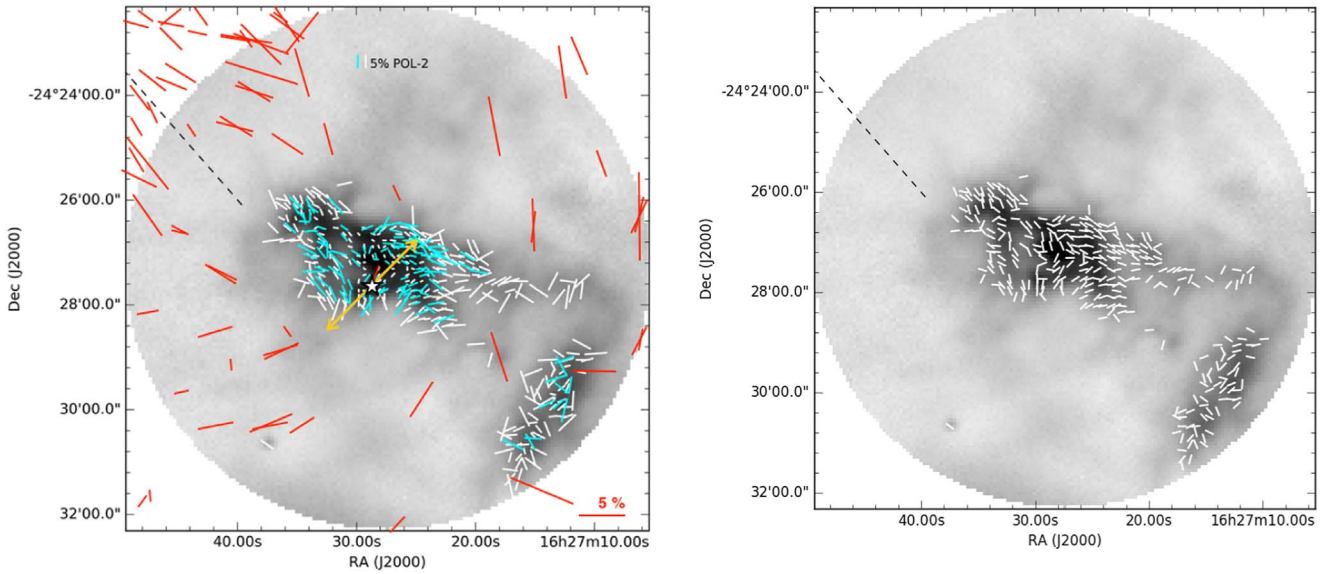
samples peaks at  $\sim 10^\circ$ – $20^\circ$ . Detailed quantitative comparison of POL-2 data with SCUPOL is difficult because of the poor number statistics and low S/N in the SCUPOL data. However, we note that the two sets of vectors are broadly consistent, and agree best in the center of the SCUPOL field. Also, note that POL-2 data have much higher sensitivity than SCUPOL, thus providing more reliable polarization angles.

## 4. Discussion

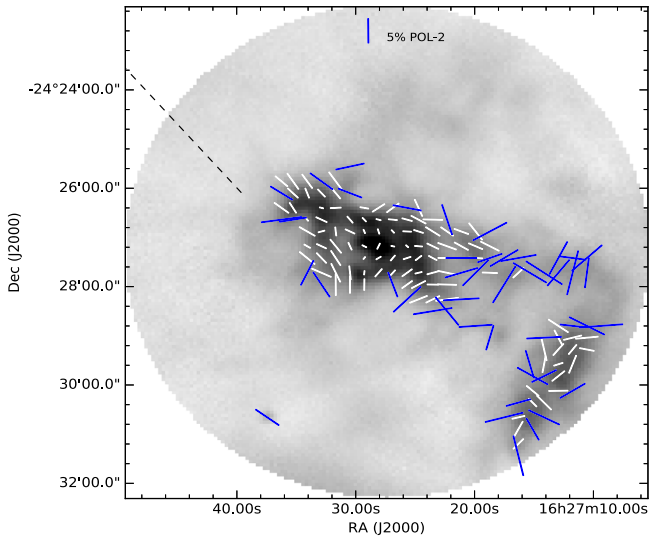
### 4.1. The Magnetic Field in NIR Wavelengths

Kwon et al. (2015) presented the  $B$ -field morphology across Ophiuchus, measured using dust extinction in the J, H, and K bands using SIRPOL.<sup>63</sup> Near-infrared (NIR) polarimetry can probe  $B$ -fields in diffuse environments ( $A_V \sim 10$ – $20$  mag; Tamura 1999; Tamura et al. 2011). However, these data cannot probe the high-density regions where stars themselves form. To study the  $B$ -field morphology on  $\lesssim 1$  pc scales, Kwon et al. (2015) compared the NIR data with optical polarimetry data tracing  $B$ -fields on 1–10 pc scales in the Ophiuchus region (Vrba et al. 1976). Their investigation suggested that the  $B$ -field structures in the Ophiuchus clumps were distorted by the cluster formation in this region, which may have been induced by a shock compression by winds and/or radiation from the Scorpius–Centaurus association toward its west. The overall  $B$ -field observed at NIR wavelengths is found to have an orientation of  $50^\circ$  east of north (Kwon et al. 2015).

<sup>63</sup> SIRIUS camera in POLarimeter mode on the IRSF 1.4 m telescope in South Africa (Kandori et al. 2006).



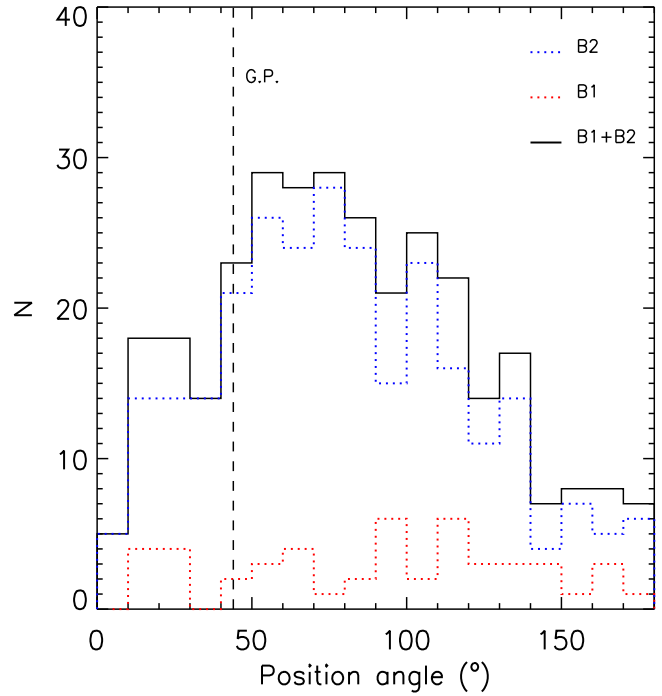
**Figure 3.** Left panel:  $B$ -field orientation in Oph-B from  $850\ \mu\text{m}$  dust polarization data. The background image shows the dust continuum map at  $850\ \mu\text{m}$  from SCUBA-2. The white and cyan color vectors correspond to data with  $PI/\delta PI > 2$  and  $PI/\delta PI > 3$ , respectively. The vectors are rotated by  $90^\circ$  to show the inferred  $B$ -field morphology. The white star symbol represents the position of IRS47, with its associated bipolar outflow shown using a double-headed arrow in yellow. A 5% polarization vector is shown for reference. The red vectors show the  $B$ -fields mapped using deep NIR observations by Kwon et al. (2015). The black dashed line shows the orientation of the Galactic Plane at the latitude of the cloud. Right panel: same as left panel (except NIR data), but the vectors are of uniform length, rather than scaled with polarization fraction.



**Figure 4.** Smoothed  $B$ -field orientations in Oph-B, shown on the  $850\ \mu\text{m}$  dust emission map from SCUBA-2. The polarization data are averaged over 4 pixels (an effective pixel size of  $8''$  using  $2 \times 2$  binning). White vectors correspond to data with  $P < 5\%$ , while blue vectors correspond to data with  $P > 5\%$ . The vectors are selected to have  $PI/\delta PI > 2$  and are rotated by  $90^\circ$  to show the inferred  $B$ -field morphology. A reference vector with 5% polarization is given for scale. The black dashed line shows the orientation of the Galactic Plane at the latitude of the cloud.

4.2. The Magnetic Field Morphology in Sub-mm Wavelengths

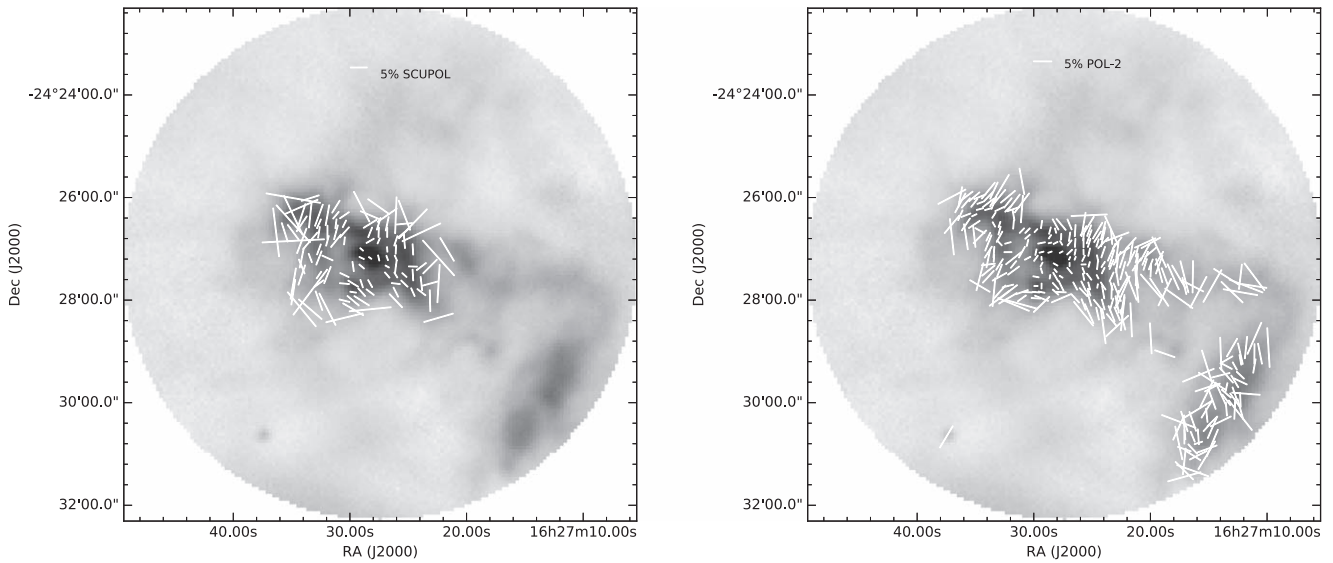
Sub-mm polarimetry probes  $B$ -fields in relatively dense ( $A_V \sim 50$  mag) regions of clouds. It can be noticed in Figure 5 that the distribution of position angles peaks at  $\theta = 50^\circ - 80^\circ$ , suggesting that the large-scale  $B$ -field is a dominant component of the cloud’s  $B$ -field geometry. There is a more clearly prevailing  $B$ -field orientation in Oph-B2, as demonstrated by the peak in the distribution shown in the histogram. The 1000–2000 au scale  $B$ -field traced by  $850\ \mu\text{m}$  polarimetry



**Figure 5.** Distributions of position angles (after  $90^\circ$  rotation) corresponding to data with  $PI/\delta PI > 2$  shown as white vectors in Figure 3. Histograms are shown for Oph-B1 in red, Oph-B2 in blue, and the entire Oph-B clump in black. The bin size is  $10^\circ$  for each distribution. The dashed line represents the orientation of the Galactic Plane.

seems to follow the structure of the two clumps, if we consider the field geometry as a whole. In Figure 3, it is clear that a  $B$ -field component with an orientation of  $50^\circ - 80^\circ$  is present in the low-density periphery of the cloud, and seems to be connecting well to the large-scale  $B$ -fields revealed at NIR wavelengths (shown as red vectors in the figure). This component is also present in Planck polarization observations





**Figure 6.** Comparison between SCUPOL (left panel) data (Matthews et al. 2009) and SCUBA-2/POL-2 (right panel) data (this work), both taken at the JCMT at  $850 \mu\text{m}$ . These are unrotated polarization vectors, not  $B$ -field orientations.

(Planck Collaboration et al. 2015). This orientation is also supported by the position angle of a single isolated detection in the southeast of the Oph-B region, which has a position angle  $\sim 60^\circ$ . The mean magnetic field direction in Oph-B2 by Gaussian fitting to the position angles is measured to be  $78^\circ \pm 40^\circ$ . The position angle of the Oph-B2 major axis is found to be  $60^\circ$  inferred by fitting an ellipse to the clump. The angular offset between the mean  $B$ -field direction and the major axis of Oph-B2 suggests that the  $B$ -field is consistent with a field that is parallel both to the long axis of Oph-B2 and to the filament-like structure in which it lies. The diffuse filamentary structure connecting the two clumps seems to have  $B$ -field lines orthogonal to it. This can be seen from the overall geometry of the  $B$ -field lines in this low-density region, which follows the clump structures. We more clearly show the smoothed  $B$ -field geometry in Figure 4. In this figure, the  $B$ -field vectors corresponding to  $P > 5\%$  and  $P < 5\%$  are plotted in blue and white, respectively. The overall  $B$ -field geometry seems somewhat disordered (see Section 4.4), with the blue vectors showing a higher degree of polarization toward the more diffuse parts of the cloud, and the white vectors showing the  $B$ -fields in the dense core regions.

### 4.3. Variation of Polarization Fraction

The left-hand panel in Figure 8 shows a map of polarization intensity toward Oph-B, with contours of Stokes  $I$  emission overlaid. The figure shows a decrease in polarization fraction in denser regions (as traced by the Stokes  $I$  emission). This depolarization has also been seen in many previous studies (e.g., Dotson 1996; Matthews & Wilson 2002; Girart et al. 2006; Tang et al. 2009a, 2009b, 2013). It is noticed that the polarization fraction is relatively higher on the edges of the cloud than in the denser core regions (see Figure 4). Plausible causes for the polarization fraction hole effect have been discussed in previous studies, such as Hull et al. (2014). They found that in low-resolution maps (i.e.,  $\sim 20''$  resolution, obtained from SHARP,<sup>64</sup> Hertz and SCUBA), depolarization resulted from unresolved structure being averaged across the

beam. However, in high-resolution maps (i.e.,  $\sim 2''$  resolution, obtained using CARMA<sup>65</sup>), depolarization persisted even after resolving the twisted plane-of-sky  $B$ -field morphologies.

Except for very few lines of sight through the densest parts of protostellar disks, (sub)millimeter-wavelength thermal dust emission is optically thin. Therefore polarization data at these wavelengths will be an integration of the emission along the entire line of sight. If the  $B$ -field changes along the line of sight (e.g., due to turbulence), the observed dust polarization will be averaged, and thus will result in a lower observed polarization fraction. Future higher-resolution ALMA observations can help trace the magnetic field structure at different scales in the cloud (e.g., Hull et al. 2013, 2014).

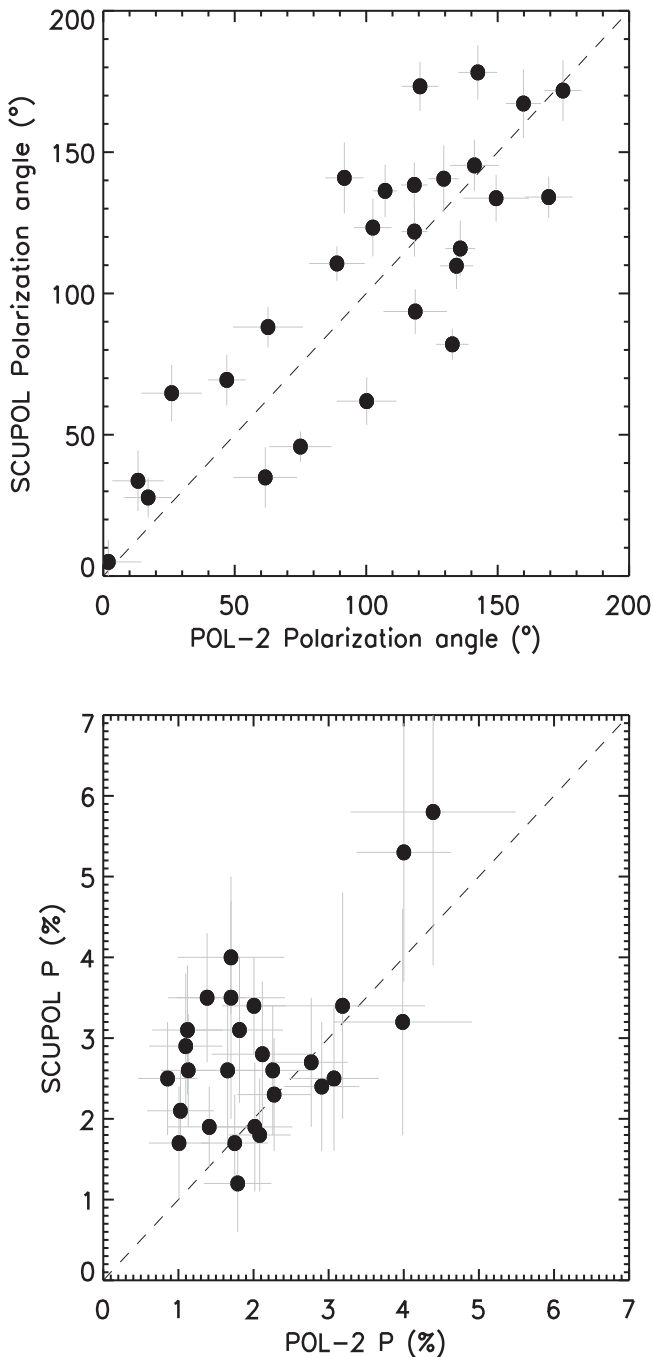
The polarization hole effect has also been explained as resulting from different grain characteristics in dense cores than in diffuse regions. Grain growth can make grains more spherical and thus harder to align with  $B$ -fields by radiative torques. Grain alignment can also be disturbed due to collisions in the higher density regions resulting in low polarization.

Alternatively the polarization hole effect could arise from a loss of photon flux to drive grain alignment in very dense regions. According to modern grain alignment theory, dust grains are aligned by radiative torques (Lazarian & Hoang 2007; Hoang & Lazarian 2008, 2009, 2016). Toward a dense region, the magnitude of these Radiative Alignment Torques (RATs) decreases due to the attenuation of the interstellar radiation. As a result, grain alignment may be suppressed in very high extinction regions, resulting in a decrease in the observed fractional polarization.

Wolf et al. (2003) found that the polarization fraction in molecular clouds follows the relation  $P \propto I^{-\alpha}$ , with  $\alpha \sim 0.5-1.2$ . They also reported a value of  $\alpha$  of  $-0.43$  toward some Bok globules. Arce et al. (1998) observed a similar behavior in the dust extinction polarization of background starlight. Falceta-Gonçalves et al. (2008) performed a statistical investigation of the previous studies in order to understand this trend. They produced synthetic polarimetric maps and found an anti-correlation between polarization degree and column

<sup>64</sup> The SHARC-II polarimeter for CSO.

<sup>65</sup> Combined Array for Research in Millimeter Astronomy.



**Figure 7.** Comparison of polarization angle (top) and fraction (bottom) values obtained toward Oph-B2 by SCUPOL (Matthews et al. 2009) and SCUBA-2/POL-2 (this work). Dashed lines show a one-to-one relation (see text for details).

density with a value of  $\alpha$  of  $-0.5$  (assuming perfect grain alignment with no dependence on extinction), due to an increase in the dispersion of the polarization angle along the line of sight in the denser regions. Cho & Lazarian (2005), in their study of grain alignment, claimed that the power-law index is sensitive to the grain size distribution.

Measured values of  $\alpha$  can vary in different cases. For example, Matthews & Wilson (2000) found a relation of  $P \propto I^{-0.7}$  toward OMC-3, but  $P \propto I^{-0.8}$  in the Barnard 1 cloud. In the case of dense star-forming cores, Henning et al. (2001) measured  $P \propto I^{-0.6}$ , whereas Lai et al. (2002) found

$P \propto I^{-0.8}$ , and Crutcher (2004) obtained a slightly steeper value of  $\alpha$ , with  $P \propto I^{-1.2}$ . The relationship between polarization fraction and intensity may be different in dense cores. Hull et al. (2017a) found a relatively high polarization fraction near the total intensity peak of Serpens SMM1 using ALMA data. This higher degree of polarization may possibly be due to additional radiative torques from its young central protostar.

The right-hand panel of Figure 8 compares degree of polarization with  $850 \mu\text{m}$  dust emission intensity in Oph-B. We measured the power-law slope in the distribution using a least-squares fit, and found  $\alpha \approx -0.9$ . A similar value has been reported by Alves et al. (2014) toward starless cores. Most recent ALMA observations at  $0''.26$  resolution toward W51 e2, e8, north show a slope of  $-0.84$  to  $-1.02$  and SMA observations at  $2''$  resolution toward the same source in e2/e8, and north reported a slope of  $-0.9$  and  $-0.86$ , respectively (Koch et al. 2018). Higher-resolution polarization observations (i.e., using ALMA) are needed to probe the magnetic fields toward the very centers of the embedded cores in Oph-B, on the scales  $< 1000$  au.

#### 4.4. The Davis–Chandrasekhar–Fermi Analysis for $B$ -field Strength

We use the Davis–Chandrasekhar–Fermi (DCF; Davis 1951; Chandrasekhar & Fermi 1953) method to estimate the plane-of-sky magnetic field strength in Oph-B. This method assumes that the underlying magnetic field geometry of the region under consideration is uniform, and hence that the observed dispersion in position angle is a measure of the distortion in the field geometry caused by turbulence, and also that the distribution of vectors about the mean field direction is approximately Gaussian, and thus well-characterized by its standard deviation.

The DCF method estimates the plane-of-sky  $B$ -field ( $B_{\text{pos}}$ ) using the formula

$$B_{\text{pos}} = Q \sqrt{4\pi\rho} \frac{\sigma_v}{\sigma_\theta}, \quad (5)$$

where  $\rho$  is the gas density,  $\sigma_v$  is the 1D non-thermal velocity dispersion of the gas,  $\sigma_\theta$  is the dispersion in polarization angle, and  $Q$  is a factor of order unity that accounts for variations in the  $B$ -field on scales smaller than the beam.

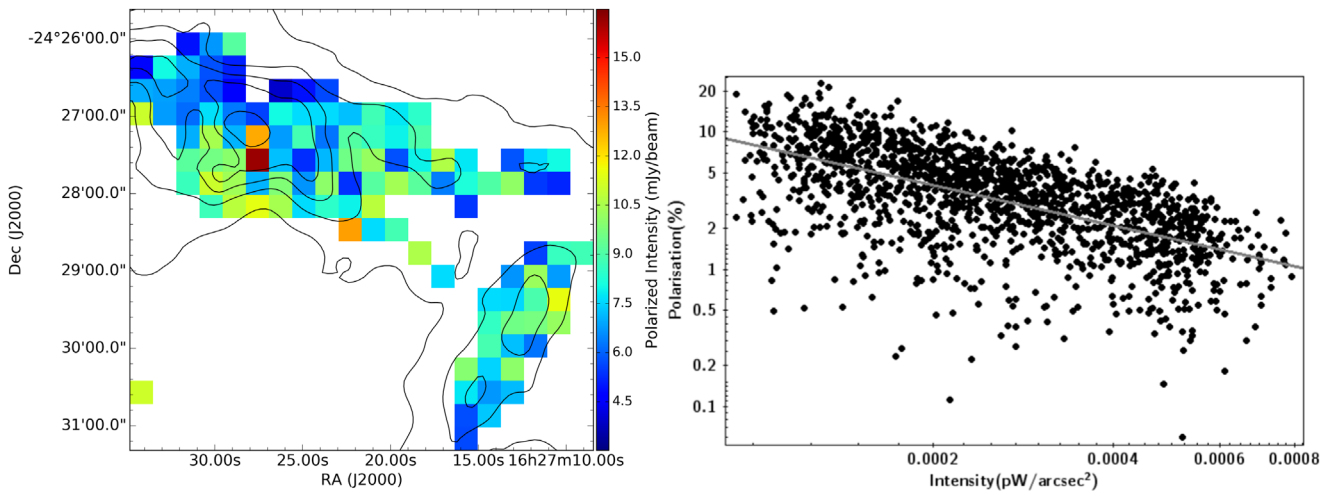
Crutcher (2004) formulated this as

$$B_{\text{pos}} \approx 9.3 \sqrt{n(\text{H}_2)} \frac{\Delta v}{\sigma_\theta} \mu\text{G}, \quad (6)$$

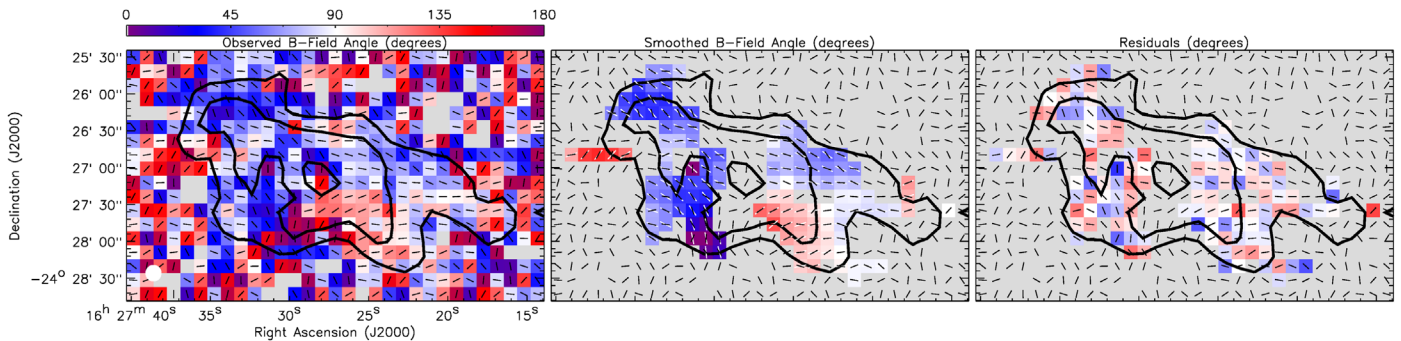
where  $n(\text{H}_2)$  is the number density of molecular hydrogen in  $\text{cm}^{-3}$ ,  $\Delta v = \sigma_v \sqrt{8 \ln 2}$  is in  $\text{km s}^{-1}$ ,  $\sigma_\theta$  is in degrees, and  $Q$  is taken to be 0.5 (cf. Ostriker et al. 2001).

##### 4.4.1. Number Density

Motte et al. (1998) infer a volume density in Oph-B of  $n(\text{H}_2) = 3.2 \times 10^6 \text{ cm}^{-3}$  from 1.3 mm dust continuum measurements. We adopt this value throughout the subsequent analysis. These authors estimate that their measured column densities are accurate to within a factor  $\lesssim 2$ , so we conservatively adopt a fractional uncertainty of  $\pm 50\%$  on their stated volume density.



**Figure 8.** Left panel shows the map of polarization intensity with dust continuum contour overlaid. The degree of polarization versus total intensity plot is shown in the right panel after selecting the sample using  $I > 0$ ,  $\frac{I}{\delta I} > 20$ , and  $P > 0$ . The gray line shows the linear fitting in the data.



**Figure 9.** Left panel: the observed magnetic field angles in Oph-B2. Middle panel: the magnetic field angle map smoothed with a  $3 \times 3$  pixel boxcar filter. Right panel: the residual map from subtracting the smoothed map from the observations. Contours show Stokes  $I$  emission at 1.0, 2.0, and 3.0  $\text{mJy arcsec}^{-2}$ .

#### 4.4.2. Velocity Dispersion

André et al. (2007) observed Oph-B in the  $\text{N}_2\text{H}^+ J = 1 \rightarrow 0$  line with beam size  $26''$  using IRAM 30 m telescope. Pattle et al. (2015) showed that there is good agreement between core masses derived using these  $\text{N}_2\text{H}^+$  data and core masses derived using SCUBA-2  $850 \mu\text{m}$  continuum observations of Oph-B (see their Figure 8). Hence we assume that  $\text{N}_2\text{H}^+$  traces the same density region as our SCUBA-2/POL-2 observations do, and so that the non-thermal velocity dispersions measured by André et al. (2007) are representative of the non-thermal motions perturbing the magnetic field, which we observe in Oph-B.

André et al. (2007) list non-thermal velocity dispersions for 19 of the cores in Oph-B2 identified by Motte et al. (1998) in their Table 4 (note that we use the non-background-subtracted values in order to have the best possible correspondence with the integrated dust column). Averaging these velocity dispersions, we estimate  $\sigma_v = 0.24 \pm 0.08 \text{ km s}^{-1}$  (i.e.,  $\Delta v = 0.56 \pm 0.20 \text{ km s}^{-1}$ ), where the uncertainties give the standard deviation on the mean.

#### 4.4.3. Angular Dispersion

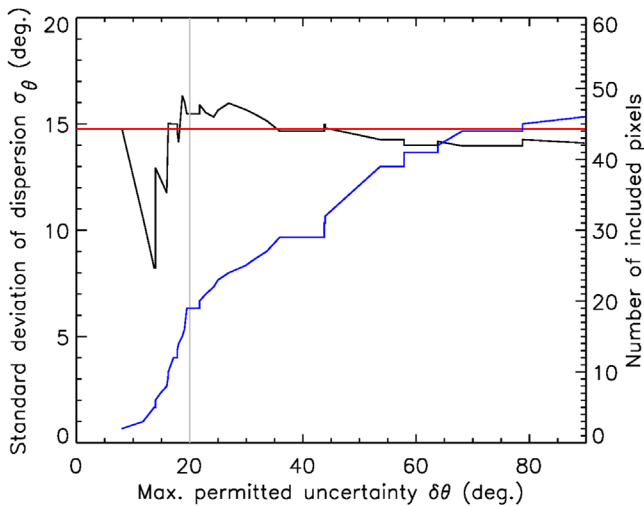
To ensure that we are measuring the dispersion of statistically independent pixels, and to maximize the fraction of the map with high S/N measurements, we gridded the vector map to  $12''$  (approximately beam-sized) pixels. We performed the DCF analysis on Oph-B2 only, because the field in Oph-B1 is quite disordered, as seen in Figure 5, with no clear peak in

the distribution of angles. The DCF method thus cannot be meaningfully applied in this region.

There are many methods of estimating the angular dispersion in polarization maps. The polarization vectors in Oph-B2 show a relatively ordered morphology with some scatter along the edges. Naively taking the standard deviation of all pixels in Oph-B2 with  $\text{PI}/\delta\text{PI} > 2$ , we measure  $\sigma_\theta = 42^\circ$ . This value is significantly larger than the maximum value at which the standard DCF method can be safely applied ( $\lesssim 25^\circ$ ; Heitsch et al. 2001). This large standard deviation includes both the scatter in polarization vectors and any underlying variations in the ordered field.

Instead, we followed the technique to measure the dispersion of polarization angle presented by Pattle et al. (2017). They used an unsharp-masking method to remove the underlying ordered magnetic field structure from their observations. First, they estimated the general magnetic field structure by smoothing their map of magnetic field with a  $3 \times 3$  pixel boxcar filter. Second, they subtracted the smoothed map from the observed magnetic field angle map. The resulting residual map was taken to represent the non-ordered component of the magnetic field. Finally, they used the standard deviation in the residual map to represent the standard deviation in the magnetic field angle,  $\sigma_\theta$ . We refer readers to Pattle et al. (2017) for a detailed description of the methodology.

We applied the Pattle et al. (2017) method to our Oph-B2 data. Figure 9 shows the observed projected magnetic field (left



**Figure 10.** Black line: the standard deviation ( $\sigma_\theta$ ) of the distribution of the magnetic field angles in Oph-B2 about the mean magnetic field direction as a cumulative function of maximum uncertainty in the  $3 \times 3$  pixel smoothing box (left-hand axis). Red line: the mean value of standard deviation, measured over the distribution to the right of the gray line (left-hand axis). Blue line: the number of pixels included in the cumulative distribution (right-hand axis).

panel), the inferred ordered component to the field (middle panel), and the residual map (right panel). We restrict our analysis to pixels where  $P/\delta P > 3$  (although the S/N may be lower in the surrounding pixels, which contribute to the mean field estimate), and exclude any pixel where the maximum difference in angle between pixels within the boxcar filter is  $\geq 90^\circ$ , on the assumption that the field in the vicinity of that pixel is not sufficiently uniform for the smoothing function to be valid. We measure a standard deviation in the residual map of  $\sigma_\theta = 14.8 \pm 0.7$ , as shown in Figure 10.

The method proposed by Pattle et al. (2017) was developed for data with higher S/N than we achieve in Oph-B. We tested the accuracy of this estimate of  $\sigma_\theta$  by measuring the standard deviation in angle in small regions of Oph-B2 that do not show significant ordered variation across them. We measured  $\sigma_\theta = 12.8$  in the northwest of Oph-B2,  $\sigma_\theta = 16.6$  in the northeast, and  $\sigma_\theta = 16.5$  in the southwest. This suggests that the dispersion in angle that we measure using the unsharp-masking method is representative of the true dispersion in angle possibly due to non-thermal motions in Oph-B2, and so we adopt  $\sigma_\theta \sim 15^\circ$  going forward.

#### 4.4.4. Magnetic Field Strength

Using our values of  $n(\text{H}_2) = (3.2 \pm 1.6) \times 10^6 \text{ cm}^{-3}$ ,  $\Delta v = 0.56 \pm 0.20 \text{ km s}^{-1}$ , and  $\sigma_\theta = 14.8 \pm 0.7$ , we estimate a plane-of-sky magnetic field strength in Oph-B2 of  $630 \mu\text{G}$ . Combining our uncertainties using the relation

$$\frac{\delta B_{\text{pos}}}{B_{\text{pos}}} = \frac{1}{2} \frac{\delta n(\text{H}_2)}{n(\text{H}_2)} + \frac{\delta \Delta v}{\Delta v} + \frac{\delta \sigma_\theta}{\sigma_\theta}, \quad (7)$$

where  $\delta n(\text{H}_2)$  represents the uncertainty on  $n(\text{H}_2)$  and so forth, we find a total fractional uncertainty of  $\delta B_{\text{pos}}/B_{\text{pos}} \approx 66\%$ . We note that a significant fraction of the contribution to this uncertainty is systematic, and thus that our result of  $B_{\text{sky}} = 630 \pm 410 \mu\text{G}$  suggests that the magnetic field strength in Oph-B2 is possibly within the range 200–1000  $\mu\text{G}$ .

#### 4.4.5. Small-angle Approximation

The DCF method fails at large angular dispersions, in part due to the failure of the small-angle approximation, as discussed in detail by Heitsch et al. (2001) and Falceta-Gonçalves et al. (2008). For our measured value of angular dispersion,  $\sigma_\theta = 15^\circ = 0.262$  radians,  $\tan(\sigma_\theta) = 0.268$  (i.e.,  $\tan(\sigma_\theta)/\sigma_\theta = 1.02$ ). We are in the regime in which the small-angle approximation holds well, and so the standard DCF method is an appropriate choice here.

#### 4.4.6. Line of Sight Turbulence

Many authors have discussed the appropriate value of the  $Q$  parameter in the DCF equation to properly account for variations if the magnetic field along the line of sight or field on scales smaller than the beam (e.g., Heitsch et al. 2001; Ostriker et al. 2001; Houde et al. 2009; Cho & Yoo 2016). Modeling suggests that the DCF method will overestimate the magnetic field strength by a factor of  $\sqrt{N}$ , where  $N$  is the number of turbulent eddies along the line of sight (Heitsch et al. 2001). Cho & Yoo (2016) propose that the number of turbulent eddies can be estimated using the relation

$$\frac{1}{\sqrt{N}} \sim \frac{\sigma_{v_c}}{\sigma_v}, \quad (8)$$

where  $\sigma_{v_c}$  is the standard deviation in the centroid velocities across the observed region, and  $\sigma_v$  is the average line of sight non-thermal velocity dispersion, as previously discussed.

Using the centroid velocities listed in Table 2 of André et al. (2007), we estimate  $\sigma_{v_c} = 0.218 \text{ km s}^{-1}$ , for the same 19 cores as were used to estimate  $\sigma_v$ . Thus we find that  $\sigma_{v_c}/\sigma_v = 0.915$  (i.e.,  $N \sim 1$ ). Thus we do not expect large numbers of turbulent eddies along the line of sight, and so any over-estimation of the magnetic field in Oph-B2 due to this effect should be minimal.

#### 4.4.7. Mass-to-flux Ratio

The relative importance of gravity and  $B$ -fields to the overall stability of a region is represented by the mass-to-flux ratio,  $\lambda$  (Crutcher 2004). The stability of cores can be estimated observationally using  $\text{H}_2$  column density and  $B$ -field strength. We assume that the average field strength from the larger Oph-B2 region applies to the dense cores within the region. Motte et al. (1998) estimated the column density of Oph-B2 to be  $41 \times 10^{22} \text{ cm}^{-2}$ . Thus, its mass-to-flux stability parameter,  $\lambda$ , can be estimated using the relation given by Crutcher (2004),

$$\lambda = 7.6 \times 10^{-21} \frac{N(\text{H}_2)}{B_{\text{pos}}}, \quad (9)$$

where  $N(\text{H}_2)$  is the molecular hydrogen column density in  $\text{cm}^{-2}$  and  $B_{\text{pos}}$  is the plane-of-the sky field strength in  $\mu\text{G}$ . From our values of  $B_{\text{pos}}$  (see Section 4.4) and  $N(\text{H}_2)$ , we find  $\lambda = 4.9 \pm 3.2$ , where the error comes from the uncertainty in  $B$ -field strength measurement only. According to Crutcher (2004) this value can be overestimated by a factor of 3 due to geometric effects. Thus, the corrected value of  $\lambda$  is  $1.6 \pm 1.1$ , which is treated as the lower limit of  $\lambda$  in this work. A value of  $\lambda > 1$  indicates that the core is magnetically supercritical, suggesting that its  $B$ -field is not sufficiently strong to prevent gravitational collapse. The values of  $\lambda$  estimated toward

Oph-B2 suggest that this clump is slightly magnetically supercritical.

#### 4.5. Relative Orientation between Outflows and Magnetic Fields

Hull et al. (2014) used CARMA observations to investigate the relationship between outflow directions and  $B$ -field morphologies in a large sample of young star-forming systems. They found that their sample was consistent with random alignments between  $B$ -fields and outflows. Nevertheless, sources with low polarization fractions (similar to Oph-B2) appear to have a slight preference for misalignment between  $B$ -field directions and outflow axes, suggesting that their field lines are wrapping around the outflow due to envelope rotation.

Oph-B2 contains a prominent outflow driven by IRS47 (White et al. 2015). The plane-of-sky position angle of outflow is  $\sim 140^\circ$ , which is roughly  $50^\circ$  offset from the general orientation of the  $B$ -field in the vicinity of the outflow. The offset between the overall  $B$ -field orientation in Oph-B2 ( $\sim 80^\circ$ ) and the outflow from IRS47 is  $\sim 60^\circ$ . These results suggest that outflows are not aligned with the  $B$ -field in Oph-B on the size scales that we observe. This is consistent with models that suggest outflows are misaligned to magnetic field orientation (Hennebelle & Ciardi 2009; Joos et al. 2012; Li et al. 2013).

### 5. Summary

This paper presents the first-look results of SCUBA-2/POL-2 observations in the  $850\ \mu\text{m}$  waveband toward the Oph-B region by the BISTRO survey. In this study, we mainly focused on the role of  $B$ -fields in star formation. The main findings of this work are summarized below:

1. We found polarization fractions and angles to be consistent with previous SCUPOL data. But our POL-2 map is more sensitive than the archival data and covers a larger area of the Oph-B clump.
2. We identify distinct magnetic field morphologies for the two sub-clumps in the region, Oph-B1 and Oph-B2. Oph-B1 appears to be consistent with having a random magnetic field morphology on the scales which we observe, whereas Oph-B2 shows an underlying ordered field in its denser interior.
3. The field in Oph-B2 is relatively well-ordered and lies roughly parallel ( $\sim 18^\circ$ ) to the long axis of the clump.
4. We find that the  $B$ -field component lying on the low-density periphery of the Oph-B2 clump is parallel to the large-scale  $B$ -field orientation identified in previous studies using observations made at NIR wavelengths. The sub-mm data show detailed structure that could not be detected with the NIR data.
5. The decrease in the polarization fraction is observed in the densest regions of both Oph-B1 and B2.
6. We measured the magnetic field strength in Oph-B2 using the DCF method. Since this sub-clump has an underlying ordered field structure, we used an unsharp-masking technique to obtain an angular dispersion of  $\sim 15^\circ$ . We measured a field strength of  $630 \pm 410\ \mu\text{G}$  and a mass-to-flux ratio of  $1.6 \pm 1.1$ , suggesting that Oph-B2 is slightly magnetically supercritical.
7. We examined the relative orientation of the most prominent outflow in Oph-B2 (IRS47) and the local magnetic field. The two have orientations that are offset

by  $\sim 50^\circ$ , suggesting they are not aligned. The angular offset between the large-scale magnetic field in Oph-B2 and the outflow orientation is  $\sim 60^\circ$ , suggesting consistency with models predicting that outflows should be misaligned to magnetic field direction.

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