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# Exploring the ex-situ components within Gaia DR3 

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#### Abstract

The presence of Gaia DR3 provides a large sample of stars with complete 6D information, offering a fertile ground for the exploration of stellar objects that were accreted to the Milky Way through ancient merger events. In this study, we developed a deep learning methodology to identify ex-situ stars within the Gaia DR3 catalogue. After two phases of training, our neural network (NN) model was capable of performing binary classification of stars based on input data consisting of 3D position and velocity, as well as actions. From the target sample of 27085748 stars, our NN model managed to identify 160146 ex-situ stars. The metallicity distribution suggests that this ex-situ sample comprises multiple components but appears to be predominated by the Gaia-Sausage-Enceladus (GSE). We identified member stars of the Magellanic Clouds, Sagittarius, and 20 globular clusters throughout our examination. Furthermore, an extensive group of member stars from GSE, Thamnos, Sequoia, Helmi streams, Wukong, and Pontus were meticulously selected, constituting an ideal sample for the comprehensive study of substructures. Finally, we conducted a preliminary estimation to determine the proportions of ex-situ stars in the thin disc, thick disc, and halo, which resulted in percentages of 0.1 per cent, 1.6 per cent, and 63.2 per cent, respectively. As the vertical height from the Galactic disc and distance from the Galactic centre increased, there was a corresponding upward trend in the ex-situ fraction of the target sample.


Key words: Galaxy: evolution - Galaxy: general-Galaxy: kinematics and dynamics - Galaxy: structure.

## 1 INTRODUCTION

The hierarchical merger hypothesis (White \& Rees 1978) has long been a prominent theory in explaining the formation of galaxies. Within the framework of the Lambda cold dark matter ( $\Lambda \mathrm{CDM}$ ) cosmology and in light of current observations, a picture has been constructed in which large galaxies are formed through a process of hierarchical mergers with smaller satellite galaxies. For instance, the ongoing merger between the Milky Way (MW) and the Sagittarius dwarf spheroidal galaxy (Sgr; Ibata, Gilmore \& Irwin 1994) provides direct observational evidence of the anticipated accretion events. Furthermore, remnants of ancient accretion events have also been observed in the form of kinematic substructures and stellar streams (e.g. Zhao, Zhao \& Chen 2009; Xue et al. 2011; Zhao et al. 2015; Helmi 2020).

Facilitated by the availability of high-quality astrometric data for a large number of stars, the Gaia mission (Gaia Collaboration 2016a) has initiated a new era in the investigation of Galactic formation. One of the most notable substructures confirmed through studies based

[^0]on Gaia data (Gaia Collaboration 2016b, 2018) is the Gaia-SausageEnceladus (GSE; Belokurov et al. 2018; Haywood et al. 2018; Helmi et al. 2018; Myeong et al. 2018). It is believed to be the remnant of the last major merger event experienced by the MW, which was completed approximately 10 Gyr ago (Gallart et al. 2019; Helmi 2020). Kinematic studies of other substructures, including Sequoia (Myeong et al. 2019), Thamnos (Koppelman et al. 2019b), and the Helmi streams (Helmi et al. 1999; Koppelman et al. 2019a) were also conducted contemporaneously.

On the other hand, spectroscopic surveys such as the RAdial Velocity Experiment (RAVE; Steinmetz et al. 2006), Sloan Extension for Galactic Understanding and Exploration (SEGUE; Yanny et al. 2009), LAMOST survey (Zhao et al. 2006, 2012; Deng et al. 2012; Liu et al. 2013; Luo et al. 2015; Yan et al. 2022), Galactic Archaeology with HERMES (GALAH; De Silva et al. 2015), Apache Point Observatory Galactic Evolution Experiment (APOGEE; Majewski et al. 2017), and H3 survey (Conroy et al. 2019), provide us with insights into chemical abundances. The formation and chemical enrichment history of dwarf galaxies (DG) differs from that of the MW, and DGs of different masses exhibit distinct chemical evolution patterns. Based on chemical abundance data obtained from the H3 survey (Conroy et al. 2019), Naidu et al. (2020) identified both previously known and newly discovered structures, such as Aleph,

Arjuna, I'itoi, and Wukong. Another chemo-dynamical study that employed data from the APOGEE DR16 (Ahumada et al. 2020) revealed evidence for the existence of a metal-poor stellar structure in the inner Galaxy (Horta et al. 2021), which was previously identified as the Kraken (Kruijssen et al. 2020) and may be a signature of the proto-MW (Belokurov \& Kravtsov 2022; Rix et al. 2022). With the support of APOGEE DR17 (Accetta et al. 2022), Horta et al. (2023) further conducted a comprehensive investigation into the chemical patterns exhibited by various previously known substructures and provided an in-depth analysis of their origins. Moreover, chemically peculiar stars were examined in APOGEE data (e.g. FernándezTrincado et al. 2019, 2022), offering further insights into the merger and evolutionary history of the MW.

The incorporation of chemical abundance significantly enhances the analysis of substructures. However, there is a substantial discrepancy in sample sizes between spectroscopic and astrometric surveys. The third data release of Gaia (Gaia Collaboration 2023) provided us with astrometric data for over 1.8 billion stars, of which more than 30 million have 6D phase space information. Concurrently, the number of sources observed by a single high-resolution spectroscopic survey is limited to the order of one million (Buder et al. 2021; Accetta et al. 2022). The LAMOST survey has achieved tens of millions of spectra, leading to numerous significant findings (e.g. Li et al. 2015; Li, Tan \& Zhao 2018; Yan et al. 2018, 2021; Xing et al. 2019, 2023; Zhao \& Chen 2021). However, as a median-to-low resolution survey, the preponderance of its data is confined to information on metal abundance and a limited number of elements. Given the current circumstances, it would be highly beneficial for the community to devise a data-driven methodology that solely relies on kinematic data to attain comparable selection outcomes as those obtained through the utilization of multi-element abundance information.

Data-driven approaches have already been utilized in researches endeavor to detect merger debris within the MW. In the pursuit of uncovering the aforementioned substructures, several classic machine learning algorithms have been employed, including KNN (Fix \& Hodges 1951), GMM (Dempster, Laird \& Rubin 1977), and HDBSCAN (Campello, Moulavi \& Sander 2013). Recent efforts to identify stellar streams and kinematic substructures have also incorporated the use of machine learning algorithms. By performing clustering using the DBSCAN algorithm (Ester et al. 1996) within a 6D phase space, Borsato, Martell \& Simpson (2020) were able to detect five high-confidence streams, including one that had not been previously discovered. Yuan et al. (2020) employed a self-organizing map algorithm (Kohonen 2001; Yuan et al. 2018) within a 4D space of orbital energy and angular momentum to identify 57 dynamically tagged groups, most of which belong to previously known substructures such as the GSE. Building upon the ANODE algorithm (Nachman \& Shih 2020), Shih et al. (2022) developed a novel methodology, termed the VIA MACHINE, for the identification of cold stellar streams within data acquired from the Gaia. Through the implementation of this approach, they were able to successfully identify the presence of the GD-1 stream (Grillmair 2006). By conducting t-Distributed Stochastic Neighbor Embedding (t-SNE; Van der Maaten \& Hinton 2008) analysis in chemical space, Ortigoza-Urdaneta et al. (2023) identified seven structures, including the Splash, GSE, the high- $\alpha$ heated-disc population, $\mathrm{N}-\mathrm{C}-\mathrm{O}$ peculiar stars, inner disc-like stars, and two previously unreported structures. Unsupervised learning algorithms do not necessitate prior labeling of data. Instead, the algorithms make decisions based on the inherent distribution of the data. Nevertheless, the absence of a labeling process makes it challenging to artificially control the behavior of unsupervised algorithms and to quantify their performance. Consequently, it cannot
be guaranteed that structures with clustering properties in phase space are related to ancient merger events. These structures may also originate from the disruption of local globular clusters (GCs) in the MW. Therefore, it is important to ensure a pure ex-situ sample before using clustering algorithms for detailed analysis.

Deep learning allows computational models that are composed of multiple processing layers to learn representations of data with multiple levels of abstraction (LeCun, Bengio \& Hinton 2015), it holds considerable promise as a tool in the search for ex-situ components. Through the implementation of supervised deep learning techniques on a known data set, an optimal model can be derived for the purpose of classification. Ostdiek et al. (2020) implemented an neural network (NN) trained to identify accreted stars using only 5D kinematics information and construct a catalogue comprising 767000 accreted stars. Subsequent companion works revealed the existence of a vast prograde stellar stream in the solar vicinity, which they named Nyx (Necib et al. 2020a, b). The methodology presented by Ostdiek et al. (2020) provides an instructive paradigm for the application of supervised learning to the identification of merger debris. By pre-training the model on simulated data, it is able to classify accreted stars labeled at truth level and capture their major kinematic features. The employment of transfer learning further improved the model's ability to adapt to real data of the MW. However, the absence of radial velocity limits the classification to kinematic level, as the derivation of integrals of motion requires full 6D phase space information. With the availability of Gaia DR3 data and the inclusion of photo-astrometric distances derived by the StarHorse code (Queiroz et al. 2018; Anders et al. 2022), the sample size of stars with accurately determined full 6D phase space parameters is no longer a constraining factor. In addition, recent data releases of spectroscopic surveys have facilitated the tagging of accreted stars in a more efficient manner, thereby enabling the improvement of the second phase of network training based on previous research.

In this study, we present a deep learning methodology for the identification of $e x$-situ stars through the utilization of 6D phase space parameters and the actions. Following the completion of two training phases, our NN model demonstrated the ability to effectively identify ex-situ components comprised of DGs, GCs, and merger debris within the target sample of Gaia DR3. We conducted a preliminary analysis of the NN classified sample and dedicate it to the community in the hope that it could provide novel insights and inspiration. Further analysis, including clustering and chemical follow-up, may be performed to fully explore the potential of this sample and deepen our understanding of the MW's merger history.

This paper is organized as follows: Section 2 introduces the NN model, the training process, and the data utilized in this study. In Section 3, we present the classification results of the NN and provide an initial analysis of the ex-situ sample. Section 4 discusses particular elements involved in our algorithm and analysis. Finally, we summarize the results of this study in Section 5.

## 2 METHOD

### 2.1 Base model

As in Ostdiek et al. (2020), we constructed an NN model that was trained on synthetic data to learn fundamental kinematic features. The NN model was developed and trained using the Keras (Chollet et al. 2015) library, with TensorFlow (Abadi et al. 2015) serving as the backend framework. The input to the NN consists of 3D position and velocity data in the Cartesian coordinate system. The model
comprises five fully connected hidden layers, with the intermediate three layers containing 128 nodes and the remaining two layers containing 64 nodes. We selected the Rectified Linear Unit (ReLU; Hahnloser et al. 2000) as the activation function, and applied batch normalization prior to each activation. The output layer of the NN is composed of a single node that uses a sigmoid activation function. This function scales the output to fall within the range of $0-1$, where values closer to 1 indicate a higher probability of the star being ex-situ. We refer to this NN model as NN_FIRE in subsequent discussions.

### 2.1.1 Synthetic data

As a supervised learning algorithm, the performance of the NN is greatly influenced by the labels of the training set. To obtain truth level ex-situ labels, we train our NN with simulated data, where the accretion history of each star is traceable. The state-of-the-art cosmological hydrodynamic simulations, specifically FIRE2 (Hopkins et al. 2018; Wetzel et al. 2023), provide us a robust platform for our training endeavors. Rather than directly employing the raw data of FIRE-2 simulations, we opted to utilize a mock Gaia catalogue constructed by Sanderson et al. (2020) to facilitate the alignment between the patterns of the simulated data and the observational data from Gaia satellite.

Detailed elaboration of the synthetic catalogue is presented in Sanderson et al. (2020), and reviewed by Ostdiek et al. (2020). The mock catalogues are derived from three MW-mass galaxies from the Latte suite of FIRE-2 simulations, namely m12i, m12f, and m12m. Assuming each star particle of mass $\approx 7000 \mathrm{M}_{\odot}$ represents a single stellar population, Sanderson et al. (2020) sampled synthetic stars following the algorithm described in Sharma et al. (2011). Three different local standard of rest (LSR) were deployed in each simulated galaxy as 'solar viewpoints' to construct the synthetic catalogue. The data structure of the synthetic catalogue resembles that of the Gaia catalogue, which supplies astrometric and photometric data. As in Ostdiek et al. (2020), we utilized data from all three LSRs in the m12i, encompassing a total of 47673267 synthetic stars. We only selected stars with available error-convolved radial velocity, which covers a distance range of $0-3 \mathrm{kpc}$. The data utilized included errorconvolved right ascension (ra), declination (dec), proper motion in right ascension (pmra), proper motion in declination (pmdec), and radial velocity (radial_velocity). In addition, the true LSR-centric distance (dhel_true) was used instead of the errorconvolved parallax (parallax) to ensure compatibility with the final model. The astrometric data were then transformed into 3D position and velocity in the Galactocentric Cartesian coordinate system using the astropy package (Price-Whelan et al. 2018), as the input of the NN.

Although the synthetic catalogue is build to mock Gaia DR2, we note that it is not necessary to reproduce the exact selection function of the latest data release of Gaia. The primary objective of NN_FIRE is to acquire fundamental knowledge of stellar kinematics. The inconsistency between the distribution of training samples and target samples will be resolved in the subsequent training process, as described in Section 2.2.

To identify the $e x$-situ stars in the synthetic catalogue, we made use of a text file named star_exsitu_flag_600.txt in the m12i directory (Wetzel et al. 2023), which provides a binary flag for each star particle at redshift equals to 0 . A flag value of 1 indicates that the star particle is $e x$-situ, while a value of 0 indicates that it is insitu. Star particles are classified as 'ex-situ' if they originated at a spherical distance greater than 30 kpc comoving from the centre of
the primary galaxy (Bellardini et al. 2022). Synthetic stars generated from the same star particle are assigned the same parentid, which is the array index of the star particle. We labeled each star with the binary flag corresponding to the index in the text file, resulting in a total of 380326 ex-situ stars.

### 2.1.2 Training and evaluation

Our training process adheres to the standard workflow within the Keras framework. We employed the Adam optimizer (Kingma \& Ba 2014 ), with an initial learning rate of $10^{-3}$. A learning rate reduction mechanism was implemented, whereby if the validation loss failed to decrease for five consecutive epochs, the learning rate would be automatically reduced to half of its original value. The minimum learning rate threshold was set at $10^{-6}$. Furthermore, to avoid overfitting, the training process would be terminated if the validation loss failed to decrease for 10 epochs. However, this early stopping mechanism was not triggered during a total of 50 epochs of training.

The synthetic stars in our data set were partitioned into training, validation, and test sets at a ratio of 6:2:2. The 6D phase space data in these three data sets were normalized using a StandardScaler from the sklearn. preprocessing module (Pedregosa et al. 2011), which was fit to the training set. The partitioning process was accomplished using the sklearn.train_test_split function, during which we employed stratified sampling on the label to ensure that the proportion of ex-situ stars remained consistent across all data sets. However, ex-situ stars comprise a very small fraction of our synthetic sample, which presents a classification problem with an extreme class imbalance. One common method to address class imbalance is sample-weighted training, but determining the appropriate weight can be challenging. If the weight is set too high, it can skew the distribution of the prediction and result in poor classification performance. Conversely, if the weight is set too low, it may not effectively address the issue of imbalanced sample size. Therefore, we aim to solve this problem by choosing a proper loss function. The Focal Loss is a loss function introduced by Lin et al. (2017) to mitigate the issue of class imbalance during training in tasks such as binary classification. This function applies a modulating term to the cross-entropy loss to focus on the samples that are challenging to classify. The function is mathematically defined as

$$
F L(p)= \begin{cases}-(1-p)^{\gamma} \log (p) & \text { if } \mathrm{y}=1  \tag{1}\\ -p^{\gamma} \log (1-p) & \text { if } \mathrm{y}=0\end{cases}
$$

where p denotes the output of the classifier and $\gamma$ represents a tunable hyper parameter. In our training process, we adopted the default setting of $\gamma=2$. As the confidence in the correct class increases, the scaling factor decays to zero, automatically downweighting the contribution of samples that are easier to classify during training and rapidly focusing the model on more challenging samples.

Due to the extreme class imbalance in our classification problem, the commonly used metric of accuracy is not a suitable measure of NN classification performance. Accuracy reflects the proportion of samples that are correctly classified within the entire sample. However, in our sample, less than 1 per cent of the stars are labeled as ex-situ. As a result, a classifier that indiscriminately assigns all stars as in-situ could easily achieve an accuracy rate exceeding 99 per cent, even in the absence of any classification activity.

In order to adopt a more appropriate metric, we initially consider ex-situ stars as positive samples while in-situ stars the negative
samples. With respect to the classification results, ex-situ samples that are accurately identified by the NN are defined as true positive samples (TP). In contrast, ex-situ samples that are erroneously classified as in-situ by the NN are defined as false negative samples (FN). Correspondingly, in-situ samples that are accurately identified by the NN are regarded as true negative samples (TN), while insitu samples that are erroneously classified as ex-situ by the NN are regarded as false positive samples (FP). To evaluate the purity of $e x$-situ samples selected by the NN, we employed the metric of precision,
Precision $=\frac{\mathrm{TP}}{\mathrm{TP}+\mathrm{FP}}$.
Furthermore, the true positive rate (TPR), commonly referred to as recall, illustrates the completeness of the selection
$T P R=\frac{\mathrm{TP}}{\mathrm{TP}+\mathrm{FN}}$.
Additionally, we calculated the false positive rate (FPR) to quantify the proportion of in-situ samples that were incorrectly classified as ex-situ,
$F P R=\frac{\mathrm{FP}}{\mathrm{TN}+\mathrm{FP}}$.
Precision and recall are often considered to be conflicting measures, making it challenging for a classifier to achieve high values for both simultaneously. In the task of identifying ex-situ stars, our goal is to ensure that the ex-situ samples selected by the NN have a high degree of purity, which corresponds to a higher precision. Similarly, we also strive to minimize the FPR. However, this may result in sacrificing the completeness of the ex-situ sample, leading to a relatively lower recall.

The output of an NN consists of a series of continuous values. In order to map these values to discrete labels, it is necessary to establish a threshold. The final classification result can be expressed as
Prediction $= \begin{cases}1 & \text { if } p>\text { threshold } \\ 0 & \text { if } p \leq \text { threshold }\end{cases}$
where $p$ represents the raw output of the NN. The choice of threshold can have a significant impact on the performance of the classifier, as different thresholds will result in varying levels of TPR and FPR. By selecting a series of thresholds, a corresponding series of TPRs and FPRs can be obtained, allowing for the construction of a receiver operating characteristic (ROC) curve. According to the performance of NN_FIRE on the test set, we present the ROC curve as shown in Fig. 1. The area under the curve (AUC) is a quantitative measure of classifier performance, with values ranging from 0 to 1 . An AUC value closer to 1 indicates superior classifier performance. In this case, NN_FIRE obtained an AUC value exceeding 0.98, denoting an excellent result. As represented by the gold star, an ideal classifier would achieve a TPR of 1 and a FPR of 0 , corresponding to a precision of 100 per cent. However, as previously discussed, this is not achievable in practice and a trade-off must be made between the purity and completeness of the ex-situ sample. When the threshold was established at 0.5 , NN_FIRE had a precision of 80.1 per cent and a recall of 34.5 per cent. Elevating the threshold to 0.75 yielded a precision of 98.3 per cent, but diminished the recall to 13.4 per cent. As Nn_FIRE is not the final model, we will refrain from discussing the threshold setting in detail at present. Based on the analysis of the ROC curve, we consider that the initial training phase has imbued NN_FIRE with the ability to identify ex-situ stars from a kinematic perspective.


Figure 1. As demonstrated by the ROC curve, the NN_FIRE achieves an AUC of 0.98 on the test set. The performance of a perfect classifier, indicated by a star, has a TPR of 1 and an FPR of 0 . The dashed line represents the performance of a random guess.


Figure 2. The parallel model comprises two components: the NN_FIRE network and a sub-network that processes actions. The outputs of these two components are concatenated and subsequently processed by a series of dense layers. The weighted layers of the NN_FIRE network are frozen, as indicated by their grey coloration.

### 2.2 Parallel model

In the context of NN regression tasks, the distribution of labels within the training set warrants substantial attention. Predictions outside this range or in sparse parts of the distribution are regarded as extrapolations of the model. Unlike regression tasks where label values are continuous, our classification task has only two categories, corresponding to label values 0 and 1 . In this case, caution should still be exercised regarding extrapolation, but the focus should be on the distribution of NN input parameters. During the training phase of NN_FIRE, only sources with distances less than 3 kpc are considered. However, the target sample may encompass distances reaching tens of kiloparsecs. Moreover, even if the simulated data is derived from a synthetic Gaia survey catalogue, it may not precisely align with observational data. Therefore, we need to design a training framework to improve our model's performance on the target sample.

Instead of performing the standard workflow of transfer learning as in Ostdiek et al. (2020), we opted for a more innovative approach. Based on NN_FIRE, a parallel model was constructed as shown in Fig. 2. The training methodology employed for this model
was similar to that of the base model, but with a shorter training period. Furthermore, the StandardScaler was supplanted by a RobustScaler to facilitate improved normalization of the realistic data. All the weights of the base model are frozen during the training process and its output is concatenated with the output of a subnetwork that takes $J_{R}, J_{z}$ and $J_{\phi}$ as input, which denotes the radial, vertical, and azimuthal actions, respectively (e.g. Binney 2012; Sanders \& Binney 2016). To enhance the normalization process, the logarithm of $J_{R}$ and $J_{z}$ was taken. The concatenated output will be further processed by the rest two dense layers, finally giving the classification result. The inputs of the sub-network introduce the potential of the MW into the parallel model and transform the model from a purely kinematic-based classifier to one that is based on dynamics. In this way, the parallel structure not only solves the extrapolation problem to a certain extent, but also maximizes the utilization of all available information of the 6D parameters. We denote this model as NN_parallel in the following text.

### 2.2.1 Observational data

The third data release of Gaia (Gaia Collaboration 2023) provided radial velocity measurement for over 30 million stars, thereby presenting an ideal target sample for our study. To overcome the limitation of parallax measurement, we employed the photoastrometric distance estimated by Anders et al. (2022). We adopted stars with available radial velocity and photo-astrometric distance within Gaia DR3, and imposed certain restrictions to ensure the quality of the data, which include a renormalized unit weight error < 1.4 (e.g. Lindegren et al. 2021), a radial velocity error $<20 \mathrm{~km} \mathrm{~s}^{-1}$, and a distance error $<30$ per cent.

In order to derive robust dynamic parameters, we employ a Monte Carlo methodology. For each observational parameter, we perform 100 times of random sampling, which follows a Gaussian distribution with the observed quantity serving as the mean and the error of the observed quantity as the standard deviation. In this way, each star is represented by 100 sets of Monte Carlo observations, which are then transformed into Cartesian positions and velocities. We apply a right-hand Galactocentric Cartesian coordinate system $(x, y, z)$, where the $x$-axis is oriented from the Sun towards the Galactic centre, the $y$-axis aligns with the direction of the MW's rotation, and the $z$-axis points towards the North Galactic Pole. We assume the circular speed at the Sun as $232.8 \mathrm{~km} \mathrm{~s}^{-1}$ and the Solar Galactocentric distance as 8.2 kpc (McMillan 2017). The distance from the Sun to the Galactic plane is set at 20.8 pc (Bennett \& Bovy 2019). For the solar peculiar motion, we adopt the value reported by Schönrich, Binney \& Dehnen (2010), yielding (U, V, W) $=(11.1$, $12.24,7.25) \mathrm{km} \mathrm{s}^{-1}$. Subsequently, we calculate the corresponding actions and orbital parameters using AGAMA (Vasiliev 2019), with the assumption of McMillan (2017) potential and the application of Stäckel approximation as described by Binney (2012). The median values of the dynamic parameters are adopted as the final estimation.

It is important to note that in the region behind the Galactic centre and close to the Galactic plane, the distances of a small number of stars may be overestimated. Since Anders et al. (2022) utilized the extinction curve provided by Schlafly et al. (2016), whose fitting quality for the inner Galaxy is relatively inferior (see fig. 10 in Schlafly et al. 2016), we consider that this overestimation could potentially be attributed to extinction issues. Consequently, we remove stars in the region defined by: $x>0 \mathrm{kpc}, z>-3 \mathrm{kpc}, \mathrm{r}_{g c}$ $>6 \mathrm{kpc},-50^{\circ}<l<50^{\circ}$ and $-10^{\circ}<b<15^{\circ}$. Here, $\mathrm{r}_{g c}$ denotes the Galactocentric distance, while $l$ and $b$ represent the Galactic
longitude and Galactic latitude, respectively. Given the asymmetric distribution of the stars whose distances may be overestimated, we adopt an asymmetric criterion. Furthermore, we exclude stars with positive total energy, leaving a target sample of 27085748 stars, which is comparable in magnitude to the sample size of stars with 5D phase space parameters in Ostdiek et al. (2020).

### 2.2.2 Chemical tagging of the training samples

For the purpose of training the NN_parallel model to be applicable to the target sample, it is necessary to construct a new data set based on observational data. The primary challenge in this process is determining an appropriate method for labeling the observational data. Ostdiek et al. (2020) implemented a selection criterion, referred to as the ZM method (e.g. Herzog-Arbeitman et al. 2018), where stars with $[\mathrm{Fe} / \mathrm{H}]<-1.5$ dex and $|z|>1.5 \mathrm{kpc}$ were labeled as $e x-$ situ. Nevertheless, given that the input of our NN consists solely of dynamic parameters, it is desirable to label ex-situ and in-situ stars using a criterion that is independent of dynamics, which ensures that the sample remains dynamically unbiased. Thus the measurement of chemical abundance is considered to be the ideal criterion for the tagging of $e x$-situ components in observations, as it offers an independent perspective, and serves as a reliable indicator of stellar origin.

In pursuit of a large sample of stars with detailed chemical abundances, we cross-matched the Gaia DR3 target sample with APOGEE DR17 catalogue (Accetta et al. 2022). Adhering to the online documentation of APOGEE DR17 and Holtzman et al. (2015), we only adopted stars that satisfy the following criteria:
(i) No BAD bit in STARFLAG or ASPCAPFLAG,
(ii) $3500 \mathrm{~K}<T_{\text {eff }}<6000 \mathrm{~K}$,
(iii) $\mathrm{S} / \mathrm{N} \geq 70$,
(iv) VSCATTER $<1 \mathrm{~km} \mathrm{~s}^{-1}$,
(v) X_Fe_FLAG $==0$ and X_FE_ERR $<0.15$ dex.

Upon applying these selection parameters, a total of 254046 stars remained. Additionally, we utilized the value-added catalogue (VAC) of LAMOST DR8 presented by Li et al. (2022), which contains abundance information for 10 elements. The elemental abundances were estimated by an NN trained on APOGEE DR17 data and show high consistency with APOGEE measurements. Therefore, we further adopted 562626 stars with X_FE_ERR $<0.10$ dex in the VAC that are in common with our target sample as a supplement to the data set, resulting in a final sample size of 816672 stars.

As demonstrated in Hawkins et al. (2015), the Galactic components can be labeled using $[\mathrm{Al} / \mathrm{Fe}]$ and $[\mathrm{Mg} / \mathrm{Mn}]$. Magnesium is the first element to be affected by Type II supernovae, while manganese is produced in higher fractions compared to iron in Type Ia supernovae. Aluminium, on the other hand, is produced by Type II supernovae and is sensitive to the initial abundance of carbon and nitrogen, which originate from helium burning or asymptotic giant branch evolution. As a result, the $[\mathrm{Mg} / \mathrm{Mn}]$ and $[\mathrm{Al} / \mathrm{Fe}]$ distributions of ex-situ stars are expected to differ from those of stars formed locally in the MW due to differences in their star formation rate and chemical evolution history. Based on APOGEE DR14 (Abolfathi et al. 2018), Das, Hawkins \& Jofré (2020) introduced the [Mg/Mn][ $\mathrm{Al} / \mathrm{Fe}$ ] diagram and discovered a'blob' of accreted stars located in the region of high $[\mathrm{Mg} / \mathrm{Mn}]$ and low $[\mathrm{Al} / \mathrm{Fe}]$, which is distinct from the stellar disc. Horta et al. (2023) and Ortigoza-Urdaneta et al. (2023) also characterized numerous substructures distributed within this region, including the GSE, Sequoia, and Heracles. Furthermore, recent studies have attempted to extract the ex-situ components in


Figure 3. As depicted by the red dashed line in the $[\mathrm{Mg} / \mathrm{Mn}]-[\mathrm{Al} / \mathrm{Fe}]$ plane, our criterion divides the plane into two distinct regions. The region to the left of the line is predominantly occupied by ex-situ stars, while the region to the right is primarily composed of in-situ stars. The middle and right subplots illustrate the distribution of stars labeled as $e x$-situ in $E-J_{\Phi}$ space and velocity space, respectively. Each subplot is colour-coded to represent the number density of stars, with warmer colours indicating higher densities for $e x$-situ stars and darker colours indicating higher densities for in-situ stars.
the $[\mathrm{Mg} / \mathrm{Mn}]-[\mathrm{Al} / \mathrm{Fe}]$ diagram (e.g. Horta et al. 2021; Carrillo et al. 2024; Feltzing \& Feuillet 2023). However, most of these efforts have focused primarily on the GSE region. As noted by Hasselquist et al. (2021) and Fernandes et al. (2023), satellite galaxies of the MW such as the Large Magellanic Cloud (LMC), the Small Magellanic Cloud (SMC), and the Sgr exhibit relatively low [ $\mathrm{Mg} / \mathrm{Mn}$ ] values. With the aim of uniformly distinguishing between the ex-situ and in-situ components, it may be more effective to select a region that encompasses as many DGs and substructures as possible, rather than focusing on a specific accretion component such as the GSE.

We performed a segmental analysis of our training samples, as depicted in Fig. 3. Stars were labeled as ex-situ if they satisfy the following selection criteria:

$$
\begin{cases}{[\mathrm{Al} / \mathrm{Fe}]<0.00} & \text { if }[\mathrm{Mg} / \mathrm{Mn}] \geq 0.45  \tag{6}\\ {[\mathrm{Mg} / \mathrm{Mn}]>1.6[\mathrm{Al} / \mathrm{Fe}]+0.45} & \text { if } 0.13 \leq[\mathrm{Mg} / \mathrm{Mn}]<0.45 \\ {[\mathrm{Al} / \mathrm{Fe}]<-0.20} & \text { if }[\mathrm{Mg} / \mathrm{Mn}]<0.13\end{cases}
$$

According to these criteria, 11478 ex-situ stars were selected in our data set. This selection includes member stars of the MW's satellite galaxies and GC debris (e.g. Fernández-Trincado et al. 2022). Inevitably, it also encompasses a minor fraction of stars associated with the Galactic disc. As shown in Fig. 3, there is a group of stars with a $v_{\Phi}$ of approximately $230 \mathrm{~km} \mathrm{~s}^{-1}$, whose distribution on the $E-J_{\Phi}$ plane closely resembles that of the disc. Fernandes et al. (2023) showed that these disc stars occupy the region of low [ $\mathrm{Al} / \mathrm{Fe}$ ] and low $[\mathrm{Mg} / \mathrm{Mn}]$, which overlaps with the ex-situ region as defined by equation (6). However, a criterion stringent enough to exclude these disc stars would also fail to capture the region where the majority of DGs reside. Upon careful evaluation, we determined that this classification method is reasonable and the pollution of disc stars will not have a significant impact on the training of NN_parallel.

### 2.2.3 Model performance

Since the weights of the base model are frozen, the chemical criterion turns out to be sufficient and robust enough to support the subsequent training of the NN_parallel. During the training phase, the training samples were divided into training set, validation set, and test set as in Section 2.1.2. On the test set, NN_parallel achieved an AUC of 0.98 , which is comparable to that of the NN_FIRE. Employing a default threshold value of 0.5 , the NN
attained a precision of 81.0 per cent and a recall of 64.6 per cent, which already signified an equilibrium between purity and completeness. This threshold will be utilized in subsequent analysis. Nevertheless, it is imperative to note that the metric values should be regarded as intuitive references only, as they are not predicated on truth-level labels as explicated in Section 2.1.2.

Independent of metrics, we conducted a more comprehensive analysis of the classification results given by NN_parallel. As depicted in Fig. 4, the selection made by the parallel model does not precisely mirror the pattern of the label. The NN_parallel avoids selecting the disc stars that was incorrectly labeled as exsitu through chemical tagging. Conversely, stars selected by the NN but labeled as in-situ through chemical tagging exhibit chemical patterns and dynamic characteristics consistent with those of the thick disc and the GSE. Since the catalogue in Sanderson et al. (2020) did not contain information about $[\mathrm{Al} / \mathrm{Fe}]$ or $[\mathrm{Mn} / \mathrm{Fe}]$, we were not able to reproduce the workflow of the second training phase on the simulation. However, based on the results depicted in Fig. 4, it appears that NN_parallel may offer an even more appropriate classification result than chemical tagging. A similar situation arose in Ostdiek et al. (2020), who randomly selected 200000 synthetic stars from m12f and labeled them using the traditional ZM selection method. They then performed transfer learning on an NN pre-trained on m12i using this data set. After transfer learning, the NN achieved a precision of 41 per cent and a recall of 47 per cent when the threshold was set at 0.75 . When the threshold was raised to 0.95 , the precision increased to 59 per cent while the recall decreased to 13 percent. In comparison, the ZM selection method only achieved a precision of 50.9 per cent and a recall of 2.4 per cent. Our second training phase, like transfer learning, is designed to retain the results of the first training phase while adapting the model to observational data. The training process enabled the parallel model to overcome the limitations of both the base model and chemical tagging, resulting in improved performance. So far, the NN_parallel model is wellprepared for the classification task of the target sample.

## 3 RESULT

After two phases of training, the NN_parallel model was applied to the target sample from Gaia DR3, resulting in the identification of 160146 ex-situ stars. As illustrated in Fig. 5, the region of the stellar disc was virtually eliminated, with the majority of the ex-situ stars being diffusely distributed in the rest of the plane. Specifically,


Figure 4. Ex-situ stars selected by the NN_parallel model in the test set. The left panel displays the distribution in the [ $\mathrm{Mg} / \mathrm{Mn}]-[\mathrm{Al} / \mathrm{Fe}] \operatorname{diagram}$, with the selection criterion indicated by a grey dashed line. Samples located in the right part are considered false positive samples with respect to the chemical tagged label. The middle and right subplots show the distribution in $E-J_{\Phi}$ space and velocity space, respectively. The background of these two subplots represents the number density of ex-situ stars selected by the NN , with darker colours indicating higher densities. False positive samples are highlighted for emphasis, constituting 0.2 per cent of the samples in the test set.


Figure 5. Spatial distribution of 160146 ex -situ stars, as selected by the NN_parallel model, is presented within the Galactic coordinate system. Three prominent overdensities are identified, corresponding to the Large Magellanic Cloud (LMC), Small Magellanic Cloud (SMC), and Sgr, respectively. Furthermore, GCs can be found as smaller overdensities. A total of 20 GCs are present in our ex-situ sample.
three prominent spacial overdensities can be intuitively identified, corresponding to the LMC, SMC, and Sgr, respectively. Although not as conspicuous as DGs, smaller overdensities in the plot reveal the presence of GCs, with a total of 20 being found. The further selection for DGs and GCs is detailed in Section 3.1. Substructures, which are considered to be disrupted DGs that merged into the MW, cannot be directly identified in the coordinate space but are crucial for understanding the MW's merger history. The selected member stars from substructures including the GSE, Thamnos, Sequoia, Helmi streams, Wukong, and Pontus are presented in Section 3.2. As exhibited in Table 1, we summarize the sample size of the DGs, GCs, and substructures in the subsequent discussion.

Furthermore, we examined the metallicity distribution of the ex-situ sample utilizing the metallicity derived from SkyMapper Southern Survey (SMSS; Wolf et al. 2018; Onken et al. 2019; Huang et al. 2022) and the Stellar Abundance and Galactic Evolution Survey (SAGES; Fan et al. 2023; Huang et al. 2023). The SMSS is a photometric survey of the southern sky, while the SAGES covers the

Table 1. Sample size of the three DGs, 20 GCs , and six substructures included in Section 3.

| Structures | N stars |
| :--- | :---: |
| LMC | 9822 |
| SMC | 1899 |
| Sgr core | 3319 |
| Sgr stream | 1269 |
| GCs | 831 |
| GSE | 47521 |
| Thamnos | 5291 |
| Sequoia | 1714 |
| Helmi streams | 1655 |
| Wukong | 3140 |
| Pontus | 668 |
| Total | 77129 |



Figure 6. The left subplot presents the metallicity distribution of 62025 ex -situ stars derived from SMSS, while the middle subplot exhibits the metallicity distribution of 18413 ex-situ stars derived from SAGES. On the right, the subplot illustrates the correlation of iron abundance measurements for 1581 common stars between the SMSS and SAGES samples, with the diagonal indicating a one-to-one correlation. Despite slight system differences, the result demonstrate a good consistency.


Figure 7. Distribution of the LMC, SMC, Sgr core, Sgr stream and GCs on the celestial sphere. DGs are depicted in distinct colours, and the names of GCs are annotated near their corresponding positions. The background illustrates the number density of the ex-situ sample.
northern sky. By cross-matching our ex-situ sample with these two catalogues, we obtained 62025 and 18413 common stars, respectively. As shown in the left column of Fig. 6, the ex-situ stars exhibits a peak metal abundance around -1.4 dex. The peak metallicity of the SAGES sample is slightly higher at approximately -1.3 dex. We checked the agreement of metal abundance measurements for the two samples using 1581 common stars between them, resulting in a good consistency. However, the measurements of the SAGES samples were slightly higher than those of the SMSS samples, possibly due to differences in the filters utilized by the two surveys. Consequently, the peak metallicity of the ex-situ sample is likely to lie between -1.3 and -1.4 dex, which is consistent with the typical metallicity of GSE (e.g. Myeong et al. 2019; Amarante, Smith \& Boeche 2020; Wu et al. 2022). Taking into account both metallicity analysis and star counts as shown in Table 1, our ex-situ sample appears to be composed of multiple components but predominated by the GSE.

### 3.1 DGs and GCs

In our ex-situ sample, DGs and GCs are the most prominent structures, exhibiting a distinct clustering in coordinate space. As
shown in Table A1, we selected three DGs combining the position and velocity. Moreover, we selected the disrupted portion of the Sgr by cross-matching with Sgr stream catalogues compiled by Yang et al. (2019) and Ramos et al. (2022). They are shown here as a comparison of the Sgr core. The member stars of these four systems are plotted on the celestial sphere in Fig. 7, where the LMC, SMC, and Sgr core exhibit distinct spatial clustering, while the Sgr stream displays a more extended distribution. Thereafter, we removed the selected member stars of the DGs from our sample, which made the overdensity regions of the GCs more prominent, facilitating the subsequent selection.

For each GC, we primarily select its member stars using the centre coordinate and angular size $(\theta)$ from references listed in Table A2. We retain stars within a projected distance $\left(d_{\text {proj }}\right)$ to the GC centre of $d_{\text {proj }}<1.5 \times \theta$, where $1.5 \times \theta$ corresponds to three times the angular radii. Subsequently, we verify whether the selected GC candidates exhibit similar motion properties by examining their proper motion and radial velocity. Outliers in proper motion phase space are manually eliminated, and stars with radial velocity outside the range of $\mu_{\mathrm{rv}} \pm 3 \sigma_{\mathrm{rv}}$ are further removed. Here, $\mu_{\mathrm{rv}}$ represents the median radial velocity of a GC, and $\sigma_{\mathrm{rv}}$ is half the difference


Figure 8. The member stars of the six substructures, as selected using the criteria outlined in Table A3, are depicted in four distinct phase spaces. Each substructure is represented by a unique colour. The background and the contours illustrate the number density of the remaining 143006 ex -situ stars after the removal of DGs and GCs.
between the corresponding value of the 84th percentile and 16th percentile. We only selected GCs with more than 10 member stars. This resulted in a total of 20 GCs , all of which are classified as ex-situ in Belokurov \& Kravtsov (2023).

The metric values, such as precision and recall, provide an indication of the average classification ability of the overall sample. While the NN demonstrates a high degree of accuracy in identifying certain portions of the sample, its performance may be somewhat diminished when classifying more challenging portions. As DGs and GCs exhibit distinct spatial clustering, their identification by the NN_parallel represents a relatively minor test of its capabilities. The true potential of the NN is further demonstrated by its ability to accurately identify member stars of substructures that no longer cluster in coordinate space.

### 3.2 Substructures

Prior to the selection of substructures, we excluded the member stars of DGs and GCs as identified in Section 3.1, which was necessary to prevent any overlap with the regions where substructures reside in the phase space. During the selection process, we adopted an approach similar to that of Naidu et al. (2020), where each time member stars of a substructure was selected, we removed them from the sample to prevent their influence on the subsequent selection of other substructures. In total, 143006 ex-situ stars were included in our analysis. We then selected 59989 member stars from six substructures according to the order and criteria outlined in Table A3.

Our selection focuses on substructures including the GSE, Thamnos, Sequoia, Helmi streams, Wukong (also known as LMS-1), and Pontus. Fig. 8 displays the distribution of their member stars in four


Figure 9. Ex-situ fraction as a function of $|z|$. The region with $|z|$ below 0.6 kpc is considered as the thin disc, while the region with $|z|$ between 0.6 and 4 kpc is considered as the thick disc. The region above 4 kpc and below 10 kpc is occupied by the halo. Different regions are represented in distinct colours, and the average ex-situ fraction is indicated by horizontal dashed lines in the corresponding region. The red horizontal lines indicate the average $e x$-situ fraction of the target sample.
phase spaces. The second subplot presents an action diamond, with $J_{\text {total }}$ defined as the sum of $J_{R}, J_{z}$, and $\left|J_{\Phi}\right|$ (e.g. Myeong et al. 2019). In the third subplot, we characterize the semi-major axis as $\left(r_{\text {apo }}+r_{\text {peri }}\right) / 2$, and the eccentricity as $\left(r_{\text {apo }}-r_{\text {peri }}\right) /\left(r_{\text {apo }}+r_{\text {peri }}\right)$, where $r_{\text {apo }}$ and $r_{\text {peri }}$ represent the distances to the apocentre and pericentre, respectively. The selection criteria, as listed in Table A3, involve multidimensional linear cuts. Fig. 8 illustrates the projection of samples with high-dimensional features in four 2D spaces. As a result, the substructures may display clear linear boundaries in the corresponding plane.

### 3.3 Ex-situ percentage

In the presence of a large sample of high purity ex-situ stars, we are able to conduct a statistical analysis of the ex-situ fraction in the MW. To ensure the robustness of our analysis, we focus exclusively on target samples located in the region where $|z| \leq 10 \mathrm{kpc}$. In Fig. B1, we present a 2D ex-situ fraction map that exhibits significant layering along the $|z|$-direction. Based on the layering pattern, we divide our sample into three distinct regions, corresponding to the thin disc ( $|z|<0.6 \mathrm{kpc}$ ), the thick disc ( $0.6 \mathrm{kpc} \leq|z|<4 \mathrm{kpc}$ ), and the halo ( $4 \mathrm{kpc} \leq|z| \leq 10 \mathrm{kpc}$ ). The ex-situ samples, as identified by the NN_parallel model, were binned according to a bin size of $|z|=100 \mathrm{pc}$. We then plotted the ex-situ fraction as a function of $|z|$ in Fig. 9, where the $x$-axis represents the right boundary of each bin and the $y$-axis stands for the corresponding ex-situ fraction. Fig. 9 provides a clear illustration of how the proportion of ex-situ stars varies with the vertical height from the Galactic disc. It also shows the proportion of ex-situ stars in the three components of the MW. The average ex-situ percentages of the thin disc, the thick


Figure 10. Ex-situ fraction as a function of R is illustrated by the purple line, which shows the overall trend of the sample. Distributions for the thin disc, thick disc, and halo as defined in Fig. 9 are depicted in distinct colours. The standard error of each bin is indicated by the shaded area.
disc, and the halo are 0.1 percent, 1.6 per cent, and 63.2 per cent, ${ }^{1}$ respectively. Consistent with our previous knowledge, the thin disc is predominantly composed of in-situ stars. In the region of the thick disc, there is a substantial increase in the proportion of ex-situ stars, while on average, the thick disc is still primarily composed of in-situ stars. In contrast, the stellar halo is dominated by ex-situ stars, with this dominance increasing with $|z|$.

Fig. 10 presents an alternative perspective in the Galactocentric cylindrical coordinate system, illustrating the variation in the proportion of ex-situ stars with respect to their distance from the Galactic centre (R). The data is binned using a bin size of $\mathrm{R}=500 \mathrm{pc}$, and the $x$-axis represents the left boundary of each bin. As illustrated by the purple line in Fig. 10, the fraction of ex-situ stars is less than 2 per cent up until $\mathrm{R} \approx 12 \mathrm{kpc}$, and it significantly increases towards the outskirts of the MW. A general growth in the ex-situ fraction with increasing R can also be observed in both the thick disc and thin disc, with the trend being more pronounced in the thick disc. The ex-situ fraction in the thick disc is higher than that in the thin disc over the entire coverage of $R$, which is also indicated by the average values in Fig. 9. Within the halo, ex-situ stars hold a dominant position and the curve is mainly influenced by the spatial distribution of member stars from distinct substructures. When R is less than 5 kpc , the ex-situ fraction is relatively stable. As R extends beyond 5 kpc , the trend of the curve begins to be dominated by the GSE (e.g An \& Beers 2021; Wu et al. 2022), exhibiting an upward trajectory until $R$ approximates 12 kpc . Subsequently, the curve commences its decline, aligning with the spatial distribution trend of GSE member stars (e.g. Naidu et al. 2021).

## 4 DISCUSSION

### 4.1 Threshold setting

In Section 3, all the ex-situ stars are selected using a threshold value of 0.5 . This threshold was chosen as it maintains an optimal balance between precision and recall, as evidenced by the results in the test set. However, as the metrics are not based on truth-level label, our choice can only be taken as a reference. In practice, the setting of threshold may vary depending on the scientific goals. If the goal is to obtain a purer ex-situ sample, the threshold can be increased. On the

[^1]other hand, if a purer in-situ sample is desired, stars with small raw prediction values from the NN should be selected. Additionally, the adjustment to threshold value can reduce the impact of samples with inaccurate observations to a certain extent. By combining a relatively stringent threshold with a restriction on the error of observations, the quality of the required samples can be effectively ensured.

If we take one step further, it becomes apparent that the ultimate goal of both traditional linear segmentation methods (e.g. Nissen \& Schuster 2010) and deep learning methods is to obtain an ex-situ sample that conforms to experience and previous knowledge. The primary difference between the two methods lies in their performance. As demonstrated in Ostdiek et al. (2020), traditional methods struggle to balance precision and recall, with high recall often accompanied by low precision. Although deep learning methods are not currently capable of breaking this deadlock, they are able to achieve higher recall under the same precision level compared with traditional methods, and vice versa. Under certain threshold values, an NN model can even surpass traditional methods in both metrics. The dynamic classification of stars involves many dimensions of information, and NN excels at processing high-dimensional data. Traditional selection methods directly restrict dynamic parameters, providing clear physical meaning. However, this approach may overlook certain dimensions and fail to account for potential connections between parameters of different dimensions. The NN method, on the other hand, directly outputs a continuous value between 0 and 1 . All we need to do is determine a threshold to divide these continuous values into two categories. While adjusting the threshold may appear less physically meaningful than adjusting limits on dynamic parameters, it is important to note that the physical meaning is already reflected in the mapping established by the NN from high-dimensional dynamic parameters to output values. Adjusting the threshold is simply an intuitive and efficient process to obtain the final classification result. Moreover, the threshold is a hyperparameter that is independent of the model structure, enabling the attainment of different classification outcomes with out necessitating any modifications to the NN itself. By fine-tuning the threshold, classification results can be tailored to achieve various scientific objectives and meet specific requirements.

### 4.2 Result assessment

In contrast to simulations, it is extremely difficult to perfectly distinguish between in-situ and ex-situ stars in observations. Consequently, validating the ex-situ sample identified by NN_parallel presents a significant challenge, and we have refrained from quantitatively evaluating the performance of the NN. Nevertheless, as depicted in Fig. 5, the NN has effectively eliminated the majority of the disc stars. This is further evidenced by Fig. 8, which clearly shows that the area where the disc is located has been accurately deducted. In Section 3, we present a selection of member stars from DGs, GCs, and substructures within our ex-situ sample. The member stars are selected based on the criteria established in the literature, without the use of any clustering algorithm. This approach was chosen as a quick and convenient means to demonstrate the potential of our algorithm and to verify the reliability of our ex-situ samples. It should be noted that most of the selection methods employed are relatively conservative. As such, the selected member stars are not exhaustive, as evidenced by the remaining overdensities in Fig. 8. In forthcoming studies, clustering algorithms will be employed to facilitate more comprehensive investigations and detailed examination of each substructure. Furthermore, the origin of the substructures and potential correlations between various substructures will be investigated through simulations (e.g. Amarante et al. 2022).

In Section 3.3, we provided a comprehensive illustration of the variation in the proportion of ex-situ stars with respect to the vertical height from the Galactic disc and the distance from the Galactic centre. However, it is important to note that our results are subject to the selection effect of Gaia and the performance of the NN. As such, the result we present only reflect the analysis of our current sample. We divide our sample into three distinct parts according to the layering pattern in Fig. B1, which do not precisely align with the three components of the MW as defined canonically (Bland-Hawthorn \& Gerhard 2016). The estimation of the proportion of $e x$-situ stars in the thin disc, thick disc, and halo should be considered as preliminary. What we present is far from a definitive conclusion but rather a possibility for statistical exploration of the ex-situ components within the MW. Despite the current limitations, we believe that our study provides novel insights and serves as a catalyst for future research endeavors.

## 5 CONCLUSION

In this work, we developed a deep learning methodology to identify the ex-situ stars within Gaia DR3. Our base model, referred to as NN_FIRE, was built to learn fundamental kinematic features. The training of NN_FIRE was conducted on a synthetic Gaia catalogue derived from a MW-mass galaxy from the Latte suite of FIRE-2 simulations, specifically the m12i. The input for NN_FIRE consists of 3D position and velocity data, which are transformed from ra, dec, pmra, pmdec, radial_velocity, and dhel_true. The output of the NN is a series of continuous values ranging from 0 to 1 , which can be interpreted as the probability of a star being ex-situ. According to the truth-level label, the NN_FIRE model achieved an AUC exceeding 0.98 on the test set. At a threshold of 0.5 , we obtained a precision of 80.1 per cent and a recall of 34.5 per cent. By increasing the threshold to 0.75 , the precision escalated to as high as 98.3 per cent.

Based on NN_FIRE, a senior model noted as NN_parallel was built, which processes dynamic data derived from real observations of Gaia. During the training phase of NN_parallel, the weights of the base model were frozen. The output of the base model was concatenated with the output of a sub-network that processes $J_{R}, J_{z}$, and $J_{\Phi}$. We adopted stars with available radial velocity and photoastrometric distance estimated by Anders et al. (2022) within Gaia DR3 as our target sample and cross-matched it with the APOGEE DR17 catalogue and LAMOST DR8 VAC to build a data set with elemental abundance information for training NN_parallel. The stars were labeled using a chemical tagging method, as outlined in Section 2.2.2, which performed a segmental selection in the $[\mathrm{Mg} / \mathrm{Mn}]-[\mathrm{Al} / \mathrm{Fe}]$ plane. Although the chemical tagging method did mix some disc stars into the ex-situ sample, with the aid of the first training phase, NN_parallel overcame this problem and produced better results than selecting solely through chemical criteria.

Out of the 27085748 stars in the target sample, the NN_parallel model successfully identified 160146 ex-situ stars. According to the metallicity distribution, the ex-situ sample is composed of multiple components, with the GSE being the most prominent. A significant number of member stars from the LMC, SMC, Sgr, as well as 20 GCs were found in the ex-situ sample. Additionally, member stars from substructures including the GSE, Thamnos, Sequoia, Helmi streams, Wukong, and Pontus were also selected. The identification of these member stars not only verifies the reliability of our ex-situ sample and our algorithm but also provides a substantial sample for future research on DGs, GCs, and substructures. Moreover, we provided an intuitive illustration
demonstrating an increase in the ex-situ fraction as the vertical height from the Galactic disc and the distance from the Galactic centre increase. Finally, We conducted a preliminary estimation of the proportion of ex-situ stars in the thin disc, thick disc, and halo, yielding percentages of 0.1 per cent, 1.6 per cent, and 63.2 per cent, respectively.

This study serves as a prior work, with the primary objective of introducing our proposed algorithm and providing a reliable large sample of ex-situ stars. The findings delineated herein are the product of a preliminary analysis of our ex-situ sample. While these findings do not encompass the full scope of potential results, we posit that they effectively underscore the viability and promise of our proposed methodology. It is our aspiration that this research will bring the community fresh insights and inspiration regarding the evolution and merger history of the MW. Looking ahead, we plan to augment our research by incorporating simulations, conducting chemical abundance analyses, and employing clustering algorithms to delve deeper into our findings. We eagerly anticipate unearthing new revelations in this exciting field of study.

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## DATA AVAILABILITY

The simulated data utilized during the training phase is accessible at http://ananke.hub.yt. Astrometric and chemical abundance data can be obtained from the websites of Gaia (https://gea.esac.esa.int/archi ve/), APOGEE (https://www.sdss4.org/dr17/irspec/dr_synopsis) and LAMOST (http://www.lamost.org/dr8/v1.1/doc/vac), respectively. The result of this study will be shared upon reasonable request to the corresponding author.

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APPENDIX A: SELECTION CRITERIA

Table A1. Summary of the selection criteria we adopted to select DGs and the literature we consulted.

| Structures | Position | Criteria | References |
| :---: | :---: | :---: | :---: |
| LMC | $\alpha=80^{\circ} .893860, \delta=-69^{\circ} .756126$ | $\begin{gathered} d_{\mathrm{proj}} \leqq 12^{\circ} ; 160 \leqq \mathrm{RV} \leqq 348\left(\mathrm{~km} \mathrm{~s}^{-1}\right) ; \\ 1.2 \leqq \mathrm{PM}_{\alpha} \leqq 2.5\left(\mathrm{mas} \mathrm{yr}^{-1}\right) ;-0.7 \leqq \mathrm{PM}_{\delta} \leqq 1.5\left(\mathrm{mas} \mathrm{yr}^{-1}\right) \end{gathered}$ | (1)(2) |
| SMC | $\alpha=13^{\circ} .18687, \delta=-72^{\circ} .8286$ | $\begin{gathered} \bar{d}_{\mathrm{proj}} \leqq 8^{\circ} ; 71 \leqq \mathrm{RV} \leqq 220\left(\mathrm{~km} \mathrm{~s}^{-1}\right) ; \\ 0.2 \leqq \mathrm{PM}_{\alpha} \leqq 1.8\left(\mathrm{mas} \mathrm{yr}^{-1}\right) ;-1.8 \leqq \mathrm{PM}_{\delta} \leqq-0.8\left(\mathrm{mas} \mathrm{yr}^{-1}\right) \end{gathered}$ | (1)(2) |
| Sgr core | $\Lambda \approx 0^{\circ}, B \approx 1^{\circ} .5$ | $\begin{aligned} &\|B\|<10^{\circ} ; b<0^{\circ} ; 90 \leqq \mathrm{RV} \leqq 220\left(\mathrm{~km} \mathrm{~s}^{-1}\right) \\ &-3.5 \leqq \mathrm{PM}_{\alpha} \leqq-2\left(\mathrm{mas} \mathrm{yr}^{-1}\right) ;-2.5 \leqq \mathrm{PM}_{\delta} \leqq-0.5\left(\mathrm{mas} \mathrm{yr}^{-1}\right) \end{aligned}$ | (3)(4) |
| Sgr stream |  | Cross-match with Sgr stream catalogues | (5)(6) |

$\overline{(\Lambda, B)}$ are coordinates in the right-handed coordinate system with respect to the Sgr stream as in Cunningham et al. (2023), $d_{\text {proj }}$ represents the projected distance on the celestial sphere.
References: (1) Nidever et al. (2020); (2) Fernandes et al. (2023); (3) Vasiliev, Belokurov \& Erkal (2021); (4) Cunningham et al. (2023); (5) Yang et al. (2019); (6) Ramos et al. (2022).

Table A2. The table exhibits the properties of 20 GCs we selected, and the literature we referred to during the selection.

| Name | N stars | Centre $(\alpha, \delta)$ | Angular size (arcmin) | $\begin{gathered} \mathrm{PM}_{\alpha} \\ \left(\operatorname{mas~yr}^{-1}\right) \end{gathered}$ | $\begin{gathered} \mathrm{PM}_{\delta} \\ \left(\mathrm{mas} \mathrm{yr}^{-1}\right) \end{gathered}$ | $\begin{gathered} \mathrm{RV} \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \end{gathered}$ | References |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| NGC 3201 | 143 | $10^{\mathrm{h}} 17^{\mathrm{m}} 36.82^{\mathrm{s}},-46^{\circ} 24^{\prime} 44.9^{\prime \prime}$ | 18.2 | $7.870 \sim 8.771$ | $-2.363 \sim-1.438$ | $468.14 \sim 506.36$ | (1)(2) |
| NGC 5904 | 127 | $15^{\mathrm{h}} 18^{\mathrm{m}} 33.22^{\mathrm{s}},+02^{\circ} 04^{\prime} 51.7^{\prime \prime}$ | 17.4 | $3.786 \sim 4.483$ | $-10.284 \sim-9.503$ | $36.01 \sim 69.77$ | (1)(2) |
| NGC 5272 | 86 | $13^{\mathrm{h}} 42^{\mathrm{m}} 11.62^{\mathrm{s}},+28^{\circ} 22^{\prime} 38.2^{\prime \prime}$ | 16.2 | $-0.458 \sim 0.176$ | $-2.894 \sim-2.362$ | $-161.77 \sim-131.13$ | (1)(2) |
| NGC 6341 | 47 | $17^{\mathrm{h}} 17^{\mathrm{m}} 07.39^{\mathrm{s}},+43^{\circ} 08^{\prime} 09.4^{\prime \prime}$ | 11.2 | $-5.144 \sim-4.646$ | $-0.914 \sim-0.445$ | $-133.81 \sim-112.20$ | (1)(2) |
| NGC 2808 | 40 | $09^{\mathrm{h}} 12^{\mathrm{m}} 03.10^{\mathrm{s}},-64^{\circ} 51^{\prime} 48.6^{\prime \prime}$ | 13.8 | $0.787 \sim 1.271$ | $-0.172 \sim 0.618$ | $89.55 \sim 125.43$ | (1)(2) |
| NGC 7089 | 40 | $21^{\mathrm{h}} 33^{\mathrm{m}} 27.02^{\mathrm{s}},-00^{\circ} 49^{\prime} 23.7^{\prime \prime}$ | 12.9 | $3.141 \sim 3.708$ | $-2.480 \sim-1.906$ | $-17.50 \sim 10.02$ | (1)(2) |
| NGC 288 | 39 | $00^{\mathrm{h}} 52^{\mathrm{m}} 45.24^{\mathrm{s}},-26^{\circ} 34^{\prime} 57.4^{\prime \prime}$ | 19.2 | $4.033 \sim 4.328$ | $-5.877 \sim-5.592$ | $-51.54 \sim-36.89$ | (1)(4) |
| NGC 362 | 37 | $01^{\mathrm{h}} 03^{\mathrm{m}} 14.26^{\mathrm{s}},-70^{\circ} 50^{\prime} 55.6^{\prime \prime}$ | 12.9 | $6.426 \sim 6.963$ | $-2.944 \sim-2.373$ | $209.43 \sim 234.15$ | (1)(2) |
| NGC 1904 | 35 | $05^{\mathrm{h}} 24^{\mathrm{m}} 10.59^{\mathrm{s}},-24^{\circ} 31^{\prime} 27.3^{\prime \prime}$ | 8.7 | $2.292 \sim 2.623$ | $-1.704 \sim-1.473$ | $200.14 \sim 215.42$ | (2)(3) |
| NGC 6101 | 32 | $16^{\mathrm{h}} 25^{\mathrm{m}} 48.12^{\mathrm{s}},-72^{\circ} 12^{\prime} 07.9^{\prime \prime}$ | 10.7 | $1.655 \sim 1.867$ | $-0.360 \sim-0.159$ | $358.13 \sim 373.06$ | (1)(2) |
| NGC 1851 | 28 | $05^{\mathrm{h}} 14^{\mathrm{m}} 06.76^{\mathrm{s}},-40^{\circ} 02^{\prime} 47.6^{\prime \prime}$ | 11.0 | $2.053 \sim 2.273$ | $-0.886 \sim-0.498$ | $309.92 \sim 326.88$ | (1)(2) |
| NGC 6779 | 27 | $19^{\mathrm{h}} 16^{\mathrm{m}} 35.57^{\mathrm{s}},+30^{\circ} 11^{\prime} 00.5^{\prime \prime}$ | 7.1 | $-2.174 \sim-1.854$ | $1.431 \sim 1.759$ | $-148.02 \sim-125.18$ | (1)(2) |
| NGC 4590 | 25 | $12^{\mathrm{h}} 39^{\mathrm{m}} 27.98^{\mathrm{s}},-26^{\circ} 44^{\prime} 38.6^{\prime \prime}$ | 12.0 | $-2.802 \sim-2.582$ | $1.592 \sim 1.888$ | $-100.39 \sim-87.51$ | (1)(2) |
| NGC 5024 | 24 | $13^{\mathrm{h}} 12^{\mathrm{m}} 55.25^{\mathrm{s}},+18^{\circ} 10^{\prime} 05.4^{\prime \prime}$ | 12.6 | $-0.274 \sim-0.029$ | $-1.468 \sim-1.234$ | $-71.21 \sim-52.03$ | (1)(2) |
| NGC 5466 | 20 | $14^{\mathrm{h}} 05^{\mathrm{m}} 27.29^{\mathrm{s}},+28^{\circ} 32^{\prime} 04.0^{\prime \prime}$ | 11.0 | $-5.415 \sim-5.227$ | $-0.866 \sim-0.736$ | $104.73 \sim 112.14$ | (1)(2) |
| NGC 5286 | 18 | $13^{\mathrm{h}} 46^{\mathrm{m}} 26.81^{\mathrm{s}},-51^{\circ} 22^{\prime} 27.3^{\prime \prime}$ | 9.1 | $-0.074 \sim 0.520$ | $-0.390 \sim 0.049$ | $51.69 \sim 72.63$ | (1)(2) |
| NGC 6934 | 18 | $20^{\mathrm{h}} 34^{\mathrm{m}} 11.37^{\mathrm{s}},+07^{\circ} 24^{\prime} 16.1^{\prime \prime}$ | 5.9 | $-2.743 \sim-2.547$ | $-4.880 \sim-4.559$ | $-419.38 \sim-399.28$ | (1)(2) |
| NGC 1261 | 16 | $03^{\mathrm{h}} 12^{\mathrm{m}} 16.21^{\mathrm{s}},-55^{\circ} 12^{\prime} 58.4^{\prime \prime}$ | 6.9 | $1.538 \sim 1.721$ | $-2.210 \sim-2.010$ | $60.74 \sim 81.34$ | (1)(2) |
| NGC 2298 | 16 | $06^{\mathrm{h}} 48^{\mathrm{m}} 59.41^{\mathrm{s}},-36^{\circ} 00^{\prime} 19.1^{\prime \prime}$ | 9.6 | $3.184 \sim 3.406$ | $-2.253 \sim-2.065$ | $141.01 \sim 158.16$ | (1)(4) |
| NGC 6981 | 13 | $20^{\mathrm{h}} 53^{\mathrm{m}} 27.70^{\mathrm{s}},-12^{\circ} 32^{\prime} 14.3^{\prime \prime}$ | 5.9 | $-1.325 \sim-1.175$ | $-3.427 \sim-3.294$ | -334.19~-322.34 | (1)(2) |

In particular, the $\mathrm{PM}_{\alpha}, \mathrm{PM}_{\delta}$, and RV columns respectively indicate the minimum and maximum values of the proper motion and radial velocity of the samples obtained by the selection method introduced in Section 3.1. References: (1) Goldsbury et al. (2010); (2) Sinnott (1988); (3) Di Criscienzo et al. (2006); (4) Shao \& Li (2019).

Table A3. Summary of the selection criteria we employed to identify the substructures and the literature we utilized for reference.

| Substructures | Criteria | References |
| :---: | :---: | :---: |
| GSE | $\left\|L_{z}\right\|<0.5\left(\times 10^{3} \mathrm{kpc} \mathrm{km} \mathrm{s}^{-1}\right) ; 30<\sqrt{J_{r}}<50\left(\mathrm{kpc}^{1 / 2} \mathrm{~km}^{1 / 2} \mathrm{~s}^{-1 / 2}\right) ; r_{\text {apo }}<40(\mathrm{kpc})$ | (1)(2) |
| Thamnos | $\eta<-0.4 ;-1.8<E<-1.6\left(\times 10^{5} \mathrm{~km}^{2} \mathrm{~s}^{-2}\right)$ | (3)(4) |
| Sequoia | $-0.65<\eta<-0.4 ;-1.35<E<-1.0\left(\times 10^{5} \mathrm{~km}^{2} \mathrm{~s}^{-2}\right)$ | (3) |
| Helmi streams ${ }^{a}$ | $0.75<J_{\phi}<1.7\left(\times 10^{3} \mathrm{kpc} \mathrm{km} \mathrm{s}^{-1}\right) ; 1.6<L_{\perp}<3.2\left(\times 10^{3} \mathrm{kpc} \mathrm{km} \mathrm{s}^{-1}\right)$ | (4)(5) |
| Wukong ${ }^{\text {b }}$ | $\begin{gathered} 0<J_{\phi}<1.0\left(\times 10^{3} \mathrm{kpc} \mathrm{~km} \mathrm{~s}\right. \\ L_{y}>2.0\left(\times 10^{3} \mathrm{kpc} \mathrm{~km} \mathrm{~s}^{-1}\right) ; \\ \left(J_{z}-J_{r}\right) / J_{\mathrm{tot}}>0.3 ; 90^{\circ}<\arccos \left(L_{z} / L\right)<120^{\circ} \end{gathered}$ | (4)(6) |
| Pontus | $\begin{gathered} -1.72<E<-1.56\left(\times 10^{5} \mathrm{~km}^{2} \mathrm{~s}^{-2}\right) ;-470<J_{\phi}<5\left(\mathrm{kpc} \mathrm{~km} \mathrm{~s}^{-1}\right) \\ 245<J_{r}<725\left(\mathrm{kpc} \mathrm{~km} \mathrm{~s}^{-1}\right) ; \\ <J_{z}<545\left(\mathrm{kpc} \mathrm{~km} \mathrm{~s}^{-1}\right) ; 390<L_{\perp}<865\left(\mathrm{kpc} \mathrm{~km} \mathrm{~s}^{-1}\right) ; 0.5<e<0.8 ; 1<r_{\text {peri }}<3(\mathrm{kpc}) \\ 8<r_{\text {apo }}<13(\mathrm{kpc}) \end{gathered}$ | (7)(8) |

[^2]
## APPENDIX B: 2D EX-SITU FRACTION MAP



Figure B1. We illustrate the distribution of ex-situ fraction within the Galactocentric cylindrical coordinate system. The map is binned with a step size of 100 pc in both R and $|z|$ axis. The ex-situ fraction in each grid is visually represented through colour coding, which distinctly exhibits layering along the $|z|$-direction. According to the layering pattern, we divide our sample into three distinct regions: (1) a region characterized by an almost negligible presence of ex-situ stars ( $|z|<0.6 \mathrm{kpc}$ ); (2) a region where ex-situ stars are slightly interspersed ( $0.6 \mathrm{kpc} \leq|z|<4 \mathrm{kpc}$ ); (3) a region predominantly occupied by ex-situ stars ( $4 \mathrm{kpc} \leq|z|$ $\leq 10 \mathrm{kpc}$ ).

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[^1]:    ${ }^{1}$ The estimations for both the thin disc and the thick disc yield standard errors less than 0.01 per cent, while for the halo, the standard error is 0.02 per cent. All estimated values are rounded to one decimal places for consistency.

[^2]:    ${ }^{a}$ To avoid the inclusion of Sgr, we impose a restriction that $|B|>10^{\circ}$.
    ${ }^{b}$ We further exclude the sky areas where the distance may be overestimated as mentioned in Section 2.2.1, specifically, the region where $-50^{\circ}<l<50^{\circ}$ and $-10^{\circ}<b<15^{\circ}$.
    Circularity is represented with $\eta, L_{\perp}=\sqrt{L_{x}^{2}+L_{y}^{2}}$, and $e$ stands for eccentricity.
    References: (1) Feuillet et al. (2020); (2) Wu et al. (2022); (3) Koppelman et al. (2019b); (4) Naidu et al. (2020); (5) Koppelman et al. (2019a); (6) Limberg et al. (2023); (7) Horta et al. (2023); (8) Malhan (2022).

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